Outstanding visible light photocatalysis by nano-TiO₂ hybrids with nitrogen-doped carbon quantum dots and/or reduced graphene oxide

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Abstract:

Historically, titanium dioxide (TiO₂) has been one of the most extensively studied metal oxide photocatalyst, however, it suffers from a large bandgap and fast charge recombination. We report the use of green, rapid, single-step continuous hydrothermal flow synthesis for the preparation of TiO₂, and TiO₂ hybrids with reduced graphene oxide (rGO) and/or N-doped carbon quantum dots (NCQD) with leap step enhancement in photocatalytic activity. Using a solar light generator under ambient conditions with no extra oxygen gas added, we observed the evolution reaction of the model pollutant (methylene blue) in real time. Tailoring of the light adsorption to match of that solar spectrum, was achieved by a combination of materials of nano-TiO₂ hybrids of nitrogen-doped carbon quantum dots and graphene in its reduced form with a photocatalytic rate constant of ca. 25 x 10^{-5} s⁻¹. Using a diversity of state-of-the-art techniques including high-resolution transmission electron microscopy, transient photoluminescence, X-ray photoelectron spectroscopy and high accuracy, sophisticated hybrid density functional theory calculations we have gained substantial insight into the charge transfer and modulation of the energy band edges of anatase due to the presence of graphene or carbon dots, parameters which play a key role in improving drastically the photocatalytic efficiencies when compared to pristine titania. More importantly, we prove that a combination of features and materials display the best photocatalytic behaviour. This performance is delivered in a greener synthetic process, that not only produces photocatalytic materials with optimised properties and tailored visible light adsorption and efficiency, but also provides a path to industrialization.

Keywords: continuous hydrothermal flow synthesis, nitrogen-doped carbon dots, nano-titanium dioxide, graphene, photocatalyst, hybrid density functional

1. Introduction

The continued exponential population growth and consequential industrial expansion and demands on planetary resources have unleashed unprecedented and unsustainable pressure on the environment as reflected by global water contamination (for instance, 80 - 90% of wastewater is released untreated in Asia and Pacific regions into surface-water bodies) [1]. This is compounded by advancing climate change, and is of great concern, thus prompting urgent scientific research into renewable, effective, and sustainable solutions. In this regard photocatalysis, can readily harness freely available clean solar energy to degrade organic contaminants [2], reduce CO₂ into renewable solar hydrocarbon fuels [3], and generate hydrogen and oxygen by the splitting of water when exposed to sunlight [4].

Since the first report of their photocatalytic bleaching of dyes in 1938 [5], titanium dioxide (TiO_2) has been one of the most extensively studied metal oxide and remains currently the most feasible photocatalyst for industrial-scale applications. It does, however, feature a large band gap energy (3.2 eV for its anatase form), which remains a challenge in meeting the requirements of visible-light applications. Integration with an active catalytic surface which allows for efficient charge transfer and kinetics is additionally identified as a limiting factor in achieving outstanding photocatalytic performance.

To overcome these challenges, has led to the design of new materials that has propagated the development of hybrid (e.g., with metal co-catalysts) [4] and/or heteroatom doped nanostructures that deliver visible light-responsive photocatalysts. For example, the addition of carbonaceous material to the surface of TiO₂ can induce and enhance visible-light photocatalytic activity [6]. Graphene, the first 2D material made of a single or a few layers of sp^2 carbon atoms arranged in a hexagonal array and conjugated electronic structure, with a large surface area, can be hybridised in its reduced oxide form with TiO₂, to provide a photocatalytic system. These can be developed to offer enhanced charge separation in electron transfer processes and increases surface-adsorption of chemical substrates (via π - π interactions) including pollutant molecules [7]. An alternative approach has been found in application of the critically small carbon quantum dots (CODs) with dimensions < 10 nm. CQDs are composed of sp² graphitic carbon core with a shell surface consisting of a mixture of sp^2 and sp^3 carbon decorated entanglements with a variety of oxygen (and optional nitrogen) functionalities, with a poorly defined electronic structure. The CQDs offer a variety of advantages that include conductivity, a tuneable size dependant bandgap (1.8-3.1 eV) tailored to match solar light adsorption, that along with their small size, facilitates short charge transfer [8]. As such, when hybridised with TiO₂, energy levels are aligned forming heterojunctions (between π states of the CQDs and the conduction band of the semiconductor) resulting in enhancement of the charge lifetime separation [9].

Configuring new arrangements of carbonaceous mixtures with a catalytic support (TiO_2) can provide unique solutions to the photocatalytic challenges discussed via novel mechanisms. The sp² carbon atoms from the CQDs core and graphene in its reduced oxide form could be very efficient in storing and shuttling electrons, and the oxygen functional groups/dangling bonds can play an important role in interconnecting the solid components through hydrogen bonds and other physical interactions, whilst nitrogen functionalities can manipulate the ability of TiO₂ to harvest visible light [10,11]. These can be achieved by using a unique, versatile, continuous synthetic technology that produces consistent high quality and large quantities of nanohybrids of choice, which would also offer a solution to one of the challenges in photocatalytic materials industrialization [12]: scale-up.

The primary focus of this research was to establish a suitable, fast, and effective strategy for the discovery of highly efficient visible light photocatalysts. Herein, we report the rapid and tailored synthesis of nano-TiO₂ hybrids with nitrogen-doped carbon quantum dots and/or reduced graphene oxide (rGO), and their enhanced photocatalytic performance under solar light using methylene blue dye as model pollutant. The materials were produced in continuous mode using a hydrothermal flow process that mixed supercritical water (450 °C, 24.8 MPa) with precursor solution(s) in a reactor to give product in fractions of a second. In the perspective of a green, adaptive and economically efficient synthetic method, the continuous hydrothermal flow synthesis (CHFS) has attracted a growing interest for the production of a new generation of nanomaterials such as graphene, graphene quantum dots, carbon quantum dots and their nanocomposites, offering enhancement of their properties and expanding their area of applications [13]. Reproducibility, an assurance of real-time full control over of the reaction parameters, and delivering materials with optimised properties, are just a few of the benefits of CHFS [2,13–16]. Our research has previously applied these particular synthetic

properties to the production of different metal oxides [13], and various carbonaceous materials (reduced graphene oxide [19], carbon quantum dots [18,19], graphene quantum dots [20,21]). This research adds further to the development of continuous flow technology of photoactive related nanomaterials. A detailed analysis in charge behaviour of the CHFS synthesised materials through a series of complementary techniques including transient photoluminescence, X-ray photoelectron spectroscopy and density functional theory calculations are reported. High-resolution transmission electron microscopy was employed to study the morphology and particle size of nano-hybrids, UV-Vis spectroscopy to provide information about optical properties and bandgap determination, and X-ray powder diffraction, Raman spectroscopy lend complementary information about the compositions and crystallinity of nano-TiO₂ hybrids. Surface area values were also determined using nitrogen adsorption technique.

2. Experimental section

2.1 Chemicals: All the materials were purchased from commercial suppliers and used without further purification. methylene blue, anhydrous citric acid ($C_6H_8O_7$) and ammonia solution (NH_3 32%) were purchased from Fisher Chemicals (UK); titanium(IV)bis(ammonium lactate)dihydroxide solution ([$CH_3CH(O)CO_2NH_4$]₂Ti(OH)₂, 50 wt% in water) was purchased from Sigma Aldrich (UK). 15 M Ω , deionized water (ELGA Purelab system) was used in all experimental work.

2.2. Synthetic Methodology: All materials were synthesised using in-house constructed CHFS set-up[17] operated at 450 °C and at constant pressure 24.8 MPa. The simplified CHFS schematic is shown in **Schematic 1**. It consists of a water pre-heater connected to a temperature controller, three high pressure pumps used for continuous delivery of water and precursors feeds, a 'T' junction (facilitating the mixing of the precursor(s)/auxilary feeds), a counterflow reactor equipped with a thermocouple (where the supercritical water meets the precursors feed and the chemical process takes place), a post-reaction pipe-in-pipe water cooler, a back-pressure regulator (BPR) (that maintains a constant pressure in the system) and a collection vessel (for the collection of the final product aqueous mixture). Generally, in each experiment pre-mixed (at the T-junction) feeds of an aqueous solution of titanium(IV) bis(ammonium lactato) dihydroxide (0.2 M) and a carbon precursor (graphene oxide, 1.0 mg mL⁻¹ and/or citric acid, 70 mg mL⁻¹) and nitrogen source (ammonia, 1.0 M) were delivered via the T-junction to meet a concurrently flowing feed of sc-H₂O in the mixing zone (reactor), whereupon formation of nanocomposites occurred. Flow rates for each corresponding pump were set as Pump 1 = 20 mL min⁻¹, Pump 2 = 5 mL min⁻¹, Pump 3 = 10 mL min⁻¹. The product mixture was then cooled, evacuated from the reactor, and collected for further isolation, purification, and analysis.

2.3 Photocatalytic experiments: Solar simulator setup: In-house system consisting of a laboratory jack, a magnetic stirrer, and a quartz tungsten halogen lamp (Osram 1 kW R7s 22000 lm linear halogen lamp) mounted on a light intensity adjustable device cooled by an electric fan. Lamp performance was verified using Ocean Insight Flame Optical Spectrometer by measuring the tungsten halogen lamp irradiance spectrum between 200 and 800 nm (see the irradiance spectrum in **Fig. S1**). This confirmed that the artificial sunlight generated by the device contains 2.3% UV radiation and 97.7% visible radiation, and the UV radiation is made of 40.6% UV-A radiation (320 - 400nm), 6.2% UV-B radiation (280-320 nm), and 53.2% UV-C radiation (200-280 nm). It can thereby be concluded that the solar lamp used in photocatalytic experiments herein, and set under normal parameters, generated light very close to solar light and was sufficient in providing the activation energy for photocatalytic processes. Photocatalytic activities were also investigated under visible light using a MinisolTM LED solar simulator (LSH-7320, MKS Newport) irradiating at wavelengths ranging from 400 to 1100 nm (**Fig. S1**) and a light intensity of 50 mW/cm² (0.5 sun). Visible light intensity studies (0.25, 0.5, and 1 sun) were also performed on the best performing sample and apparent quantum (AQE) values were calculated (see electronic supplementary information).

The photocatalytic activities of titania nanocomposites were evaluated by the photodegradation of methylene blue (MB) in water using the solar-simulator(s) set-up as described above. Routinely, each photocatalyst (10.0 mg) was suspended in MB aqueous standard solution (0.02 mM, 50 mL) in a glass beaker (Corning Pyrex Griffin) equipped with a magnetic bead. The solution was stirred in the dark for 30 minutes to ensure the establishment of an adsorption-desorption equilibrium. The reaction mixture was located at a distance (27 cm)

from the light source with stirring maintained for the whole photocatalytic process period. The samples were irradiated for 180 minutes. During the irradiation process, aliquots (10 mL) were collected at 30 minutes intervals, centrifuged to remove particulates (5 mins at 4500 rpm), and analysed by UV-Vis spectrophotometry (Shimadzu UV-1800) to assess MB photodegradation by its absorbance change at 660 nm. After the UV-Vis analysis, the centrifuged material was resuspended in the aliquot and returned to the initial reaction mixture. The quantitative determination of MB was calculated using a calibrated correlation between the measured absorbance and its concentration. The direct degradation of MB was measured as a control in all photocatalytic measurements. Adsorption coefficients were also determined using Langmuir-Hinshelwood model (**Table S1**).

2.4 Equipment: Prior to analysis and testing, the samples were freeze-dried using a Heto PowderDry PL 3000 freeze-dryer. Siemens D5005 X-Ray Diffractometer operating at 30 kW with a long fine focus Cu K alpha radiation source ($\lambda = 1.54$ Å) was used to obtain X-ray powder diffraction (XRD) patterns. Kratos Axis Ultra DLD photoelectron spectrometer with a monochromatic Alka source operating at 144 W was utilised for Xray photoelectron spectroscopy (XPS) measurements. Samples were mounted using conductive carbon tape. Survey and narrow scans were performed at constant pass energies of 160 and 40 eV, respectively. The base pressure of the system was *ca*. 1×10^{-9} Torr rising to *ca*. 4×10^{-9} Torr under the analysis of these samples. It is prudent at this point to state that all calibration of the spectra have been made to the $Ti(2p_{3/2})$ peak taken to be 459.3 eV in line with that of Diebold and Madey. Whilst, for example, the NIST database has a mean value of ca. 458.5 eV, there is significant deviation (approx. 1.5 eV). Using lower values of the $Ti(2p_{3/2})$ peak (e.g. 458.5 eV), whilst giving stable Ti(2p) and O(1s) peak energies, there is an unacceptable low binding energies determined for sp² and sp³ carbon functions, equally fitting to the sp³ carbon taken to be 285 eV, results in a deviation in binding energies of ca. 0.6 eV For the Ti(2p) and O(1s) signals. Surface area analysis (SABET) of the powders were determined by the Branauer-Emmett-Teller (BET) method using a Micrometrics Gemini VII. The 5-point N₂ absorption method was used. The powders were the degassed prior to analysis at 180 °C for 2 hrs to remove any moisture using a FlowPrep 060 degasser. Double-corrected JEOL ARM200F, equipped with a cold field emission gun and set to 80 kV and 10 µA was employed to record high-resolution transmission electron microscopy (HRTEM) images. Prior to imaging, the aqueous solution of nanocomposites was deposited onto a holey carbon coated Cu-grid (400 µm). Image analysis was performed using ImageJ software. Scanning Electron Microscope (SEM) images were recorded using a Hitachi SU-70 FEG instrument. Fourier-Transform Infrared Spectroscopy (FT-IR) analysis was performed using a Shimadzu IRAffinity-1S spectrophotometer fitted with an Specac Quest ATR Accessory (diamond 4000-200 cm⁻¹). Raman spectroscopy analysis the as-synthesized materials was measured with a Horiba LabRAM HR Evolution spectrometer ($\lambda = 633$ nm). A Perkin-Elmer Lambda 950 UV–Vis spectrophotometer fitted with an integrated sphere was used to absorption measurements of the photocatalysts powders. The spectra were converted to Tauc plots and the bandgaps of the materials were calculated for indirect and direct transitions [2]. PL decays were measured using TCSPC with an FLS1000 Photoluminescence Spectrometer (Edinburgh Instruments, UK) equipped with a double emission monochromator and a high-speed PMT detector (H10720). The TiO₂ suspensions were excited at 375 nm by a pulsed diode laser (EPL-375) and the decays measured at an emission wavelength of 450 nm. The decays were fitted using reconvolution analysis with the IRF in the FLS1000 Fluoracle software.

2.5 Theoretical Calculations. DFT was used to calculate the equilibrium atomic geometry and electronic structure of bulk TiO₂, an unperturbed graphene sheet, and two-dimensional slabs of TiO₂, containing cuts of the (001), (100) and (101) surfaces with and without a layer of graphene adsorbed on the surfaces on either side of the slab. Plane-wave DFT as implemented in the **VASP** code was used [29] employing the hybrid exchange and correlation functional HSE06 [25], which is known to reproduce well the lattice parameters and main features of the energy band structure of TiO₂ [26]. The projector augmented wave [27] method was used to model the interaction between core and valence electrons (with four valence electrons for Ti and C and six for O). The total energy of the TiO₂ anatase primitive unit cell was calculated using a 600 eV plane-wave cut off and a $6 \times 6 \times 6$ Γ -centred Monkhorst-Pack [28], *k*-point mesh to sample the Brillouin zone (BZ), which resulted in energies converged within less than 0.01 eVÅ⁻¹. The plane-wave cut-off of 600 eV, combined with an $8 \times 8 \times 1$ *k*-mesh, resulted in total energy convergence within 0.5 meV per atom for the primitive

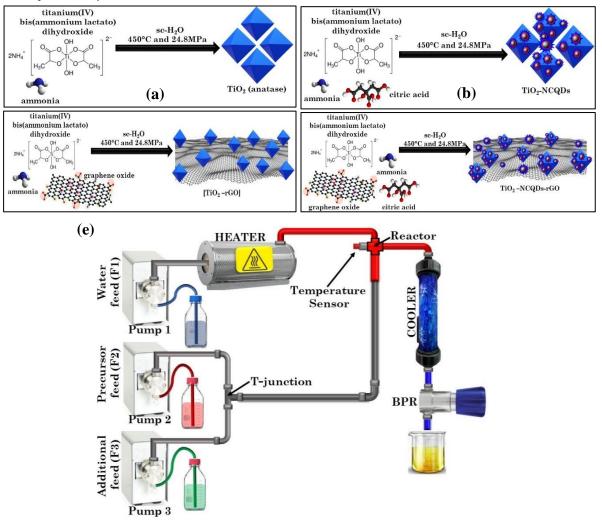
cell of the two-atom unperturbed graphene sheet. The geometry of the graphene sheet was optimised by varying the bond length about the experimental value of 1.42 Å, calculating the total energy at each length (using five data points), applying a parabolic fit to the energy versus bond length trend and finding the minimum. A vacuum gap of 15 Å was maintained between neighbouring graphene sheets in the periodic unit cell used in the simulations. The TiO₂ slabs containing the (001), (100) and (101) surfaces were produced using the surface cutting tools provided on the SAINT database [29]. In each case, a 15 Å vacuum gap was placed between periodic images of the slabs. The slab thicknesses were: 27.3 Å (001), 24.6 Å (100) and 23.5 Å (101). Geometry optimisation was performed using the same approach as that adopted for the bulk systems. The k-point meshes were chosen to maintain the density used in the graphene simulations. The slab containing the (001) surfaces was composed of 36 atoms, with a $6 \times 6 \times 1$ k-mesh used to sample the BZ; the slabs containing the (100) surfaces and (101) surfaces were both composed of 84 atoms, with a $6 \times 2 \times 1$ k-mesh used to sample the BZ. After optimisation, graphene layers were added to these slabs, placing one layer over each of the two surfaces in each slab model. To ensure the calculations were tractable, strains of over 5 % were introduced to the graphene layers, allowing them to fit within the unit cells of the slab models. For the slab containing the (001) surfaces, a 1×2 expansion of the unit cell was needed. The distance between the graphene sheet and the adjacent surface was varied between 2.5 A and 3.5 A before optimisation, to search for possible local minima. All slabs were optimised until the forces were below 0.025 eVÅ⁻¹. The macroscopically averaged electrostatic potential V [30], as a function of distance z perpendicular to the surfaces in the slab models, was computed for each slab, with and without graphene, from the calculated potential containing the Hartree and ionic contributions, using the MacroDensity package [31,32]. Combining the offset in V(z) between its value in the vacuum layer and in the interior of each slab (ΔV) with the energy of the valence band maximum of the bulk system, E_{VBM} , ionisation potentials I were calculated as $I = E_{VBM} + \Delta V$. I and electron affinity, $A = -I + E_g$, where E_g is the band gap of the bulk system, can be used to determine the band edges on an absolute scale, with the valence band (conduction band) position given by -I(A).

3. Results and Discussions

In recent years, various methods have been used to improve the photocatalytic properties of titanium dioxidecarbonaceous nanocomposites by altering or extending the TiO₂ anatase photoactivity for radiation harvesting from UV to visible [33]. Whilst established synthetic processes are promising, they still suffer drawbacks such as high energy and time consumption demands, costly precursors, or experimental setup, batch to batch quality discrepancies, and the inability to scale-up the process from the laboratory to industrial production. The continuous hydrothermal flow synthesis approach reported herein addresses all the shortcomings of traditional methods due to the systems inherent precision control of reaction parameters, that not only allows real time tunability of the material production process, but also reduced the time required to assess the best synthetic experimental conditions considerably. CHFS has previously been applied to TiO₂ synthesis [34] the uniqueness of this work is that we simultaneously and homogeneously generate titanium dioxide anatase nanoparticles, NCQDs, and reduce graphene oxide. The experimental details are shown in **Fig. 1 a-d**.

In general, pre-mixed (at the T-junction) feeds of titanium(IV) bis(ammonium lactate) dihydroxide and a carbon precursor (graphene oxide and/or citric acid) were delivered to meet concurrently with a flowing feed of sc-H₂O (450 °C and a pressure of 24.8 MPa) in the mixing zone (reactor) for nanocomposite formation. After cooling, the product mixture was collected for isolation and purification. The as-produced hybrid structures did not require any post-thermal processing; instead, the nanomaterials were characterized and used for further analysis and photocatalytic testing as described in the following.

Scheme 1: Continuous hydrothermal flow synthesis (CHFS) process detailing the synthetic protocols for: (a) -(d) production of TiO₂ and its nanocomposites with reduced graphene oxide (rGO) and/or N-doped carbon quantum dots (NCQD); (e) simplified representation of the CHFS process.



3.1 Characterisation of the photocatalysts materials

High-resolution transmission electron microscopy (HRTEM) reveals the morphological architecture of the TiO₂ and TiO₂-carbonaceous nanocomposites as shown in **Fig. 1**. The TiO₂ mean particle size is ~ 5.0 nm in diameter with a lattice spacing of *ca*. 0.33 nm (**Fig. 1e**) and assigned to the (101) plane of anatase [35]. **Table S1** shows the average particle size values of TiO₂ photocatalysts, with no variation observed (values fall within reported standard deviation). The NCQDs exhibited a spherical morphology with a mean particle size of ~ 3.4 nm with in-plane lattice spacing of *ca*. 0.22 nm (**Fig. 1e**) and identified as a graphitic core arrangement [19] with the graphene sheets consisting of 5-8 layers. The corresponding selected area electron diffraction pattern (insets in **Fig. 1a-d**) of each sample confirms the existence of the crystalline TiO₂ anatase phase and influence due to the presence of an amorphous carbon phase in the TiO₂-carbonaceous nanocomposites. A close interaction between the TiO₂ nanoparticles and NCQDs as a solid-solid direct interface was observed, as shown in **Fig. 1e**. In the case of TiO₂-rGO (**Fig. 1c**) and TiO₂-NCQDs-rGO (**Fig. 1d**), the TiO₂ nanoparticles and NCQDs are dispersed over the 2D layer, creating extensive interfacial interactions across the entire surface of the rGO layers and interlayers. Scanning electron microscopy (SEM) images for TiO₂-NCQDs-rGO (**Fig. 1f**) supported HRTEM analysis for the layered structure of reduced graphene oxide.

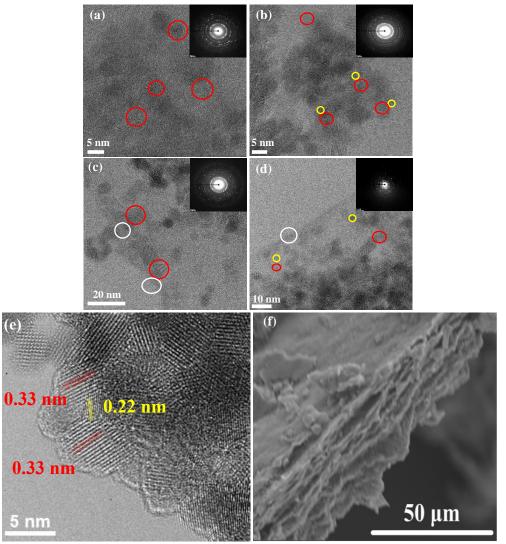


Figure 1: High-resolution transmission electron microscopy (TEM) images of: (a) TiO₂, (b) [TiO₂-NCQDs], (c) [TiO₂-rGO], (d) [TiO₂-NCQDs-rGO)]. Each TEM image has an inset of a corresponding selected area electron diffraction pattern. (e) lattice fringes of TiO₂ and NCQDs in TiO₂-NCQDs, and (f) Scanning electron microscopy image of TiO₂-NCQDs-rGO showing reduced graphene oxide (rGO) layers. Sample labelling key: yellow -NCQDs, red – TiO₂, white – rGO.

The small particle sizes observed by HRTEM correlated with the large surface area of the as-synthesised CHFS materials measurements by SA_{BET} for the pristine titania and its hybrids (compiled in **Table S1**) ranging from 232 m² g⁻¹ for TiO₂ to 253 m² g⁻¹ for TiO₂ –NCQDs-rGO nanocomposite. Such high specific area values would be expected to increase the activity of the nano-hybrids through increased adsorption of pollutant molecules (model methylene blue dye) and irradiation photons. Of course, other variables, including structural and surface properties of the photocatalysts, must be considered when comparing activities of different materials.

For further investigation of the structures and composition of TiO₂ photocatalysts and its hybrids, Raman spectroscopy and XRD analysis were undertaken (**Fig. 2**). Characteristic Raman vibrations for the CHFS-synthesised samples of TiO₂ anatase presented an E_g band located at 146 cm⁻¹ and 641 cm⁻¹, a B_{1g} band located at 398 cm⁻¹, and an A_{1g} band positioned 513 cm⁻¹, all in good agreement with the literature [36]. Adding a carbonic phase (reduced graphene oxide) to the material, not only were the expected appearance of the D (*ca.* 1345 cm⁻¹) and G (*ca.* 1595 cm⁻¹) vibrational bands observed, but these bands for all rGO containing hybrids were noticeably blue shifted (D-band) and red shifted (G band). It should also be noted that the standard E_g vibrational mode at 144 cm⁻¹ was also red-shifted for TiO₂-rGO (150 cm⁻¹), and TiO₂-NCQDs-rGO (152 cm⁻¹) respectively. This red shift could be assigned to the charge transfer process occurring between the TiO₂ nanoparticles and its hybridised carbonaceous materials [37]. This would be in line with the XPS, and modelling data discussed later. Interestingly, the D and G vibrational modes are missing for TiO₂-NCQDs,

most likely due to a too low NCQDs content in the samples for Raman analysis, but nanocomposites NCQDs are accounted for in the XPS analysis later.

Lending further insight and complementary support for our findings, XRD analysis (**Fig. 2b**) showed characteristic diffraction peaks for TiO₂ anatase (JCPDS 21-1272) in all cases. The diffraction peak at $2\theta = 12.2^{\circ}$ indicated the presence of the carbonic phase [002] and was observed for all samples. Critically, interactions between TiO₂ and NCQD and/or rGO are observed by X-ray photoelectron spectroscopy (XPS) and are discussed next.

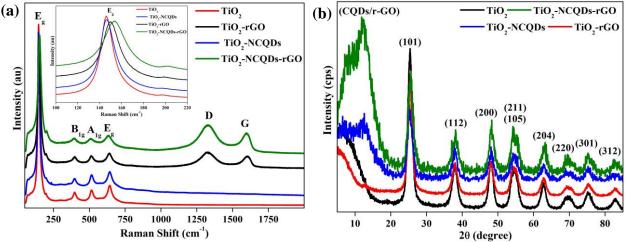


Figure 2:(a) the Raman spectra with anatase TiO₂ E_g vibrations inset and (b) the XRD patterns of CHFS-synthesised nano-TiO₂ hybrids with nitrogen-doped carbon quantum dots and/or reduced graphene oxide revealing characteristic diffraction peaks for anatase and carbonic phase.

The XPS analysis of TiO₂ nanoparticles (**Fig. 3**) reveals the characteristic Ti(2p) spin-orbit split peaks, the binding energy of the Ti(2p_{3/2}) peak recorded at 459.3 eV; the corresponding O(1s) peak is comprised of two peaks corresponding to Ti-O lattice oxygen (530.6 eV) and surface OH/CO_x (531.8 eV), a N(1s) signal at 400.7 eV is also found, consistent with fragments in the starting material for the synthesis of the nano-TiO₂. For a better understanding of the interaction between the carbon phase and oxide phase, the high-resolution XPS spectrum of carbon (1s) was analysed for all the samples, **Fig. 3a**. It is worth noting that compared to the pure TiO₂, the survey spectra of all samples exhibited higher intensity in carbon originating from graphene and/or carbon quantum dots. Addition of NCQDs to pristine TiO₂, resulted in asymmetry to the lower binding energy side of the C(1s) envelope, characteristic of the sp² graphitic core peaks at binding energies of 284.5 eV (sp³ component was measured at 285.4 eV). The difference in binding energy of the sp² component is possibly indicative of the difference in attachment of the NCQDs to the TiO₂, or the charge states of the TiO₂ and NCQDs on the surface. Additional C(1s) peaks at 286.7 and 289.4 eV were noted and ascribed to C-O/C-N and COO moieties respectively, whilst N(1s) species at 400.7 eV is again characteristic of C-N/NH_x functions.

The TiO₂-rGO complex has peaks centred at 284.5, 285.1, 286.8 and 289.4 eV, again characteristic of sp², sp³, C-O/C-N and COO functionalities as noted previously. N(1s) signals are found at 398.5 and 400.7 eV representative of C=N-C and C-N type functions. The N presence in rGO is not unexpected as it is considered a consequence of the incorporation of nitrogen from the reaction process. In the case of TiO₂-NCQDs-rGO nanomaterial, the high-resolution carbon spectra again show characteristic peaks for sp² (here the π - π * shake up structure is particularly evident) and C–N/C-O bonds, together with sp³ and the addition of C=O functions. For TiO₂-NCQDs-rGO, the N(1s) spectra, whilst showing the presence of the 400.7 eV, C-N/-NH_x peak, there is a 0.5 eV deviation in the lower binding energy species (399.0 and 398.5 eV respectively), the lower peak typically ascribed to -NH- or C=N species [38]. XPS confirms the relationship between TiO₂-NCQDs, as well as from XRD for TiO₂-rGO.

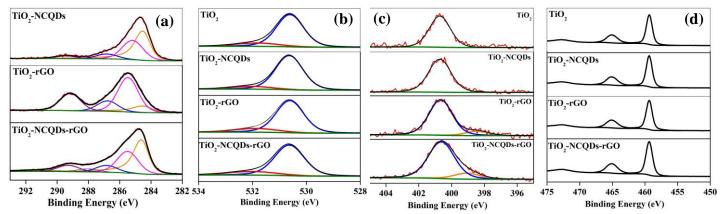


Figure 3: High resolution spectra of (a) C1s, (b) O1s, (c) N1s, and (d) Ti2p high-resolution XPS spectrum of TiO₂ and its carbonaceous nanocomposites.

FT-IR spectroscopy analysis (**Fig. 4**) reveal the distinctive overlapping stretches $(2700-3500 \text{ cm}^{-1})$ due to the presence of hydroxyl groups (-OH, -COOH) on the surface of TiO₂, and its nanocomposites, as well as protonated amino groups for TiO₂ nanocomposites of NCQDs and/or rGO, along with the expected C–H stretching vibrations (2950 cm⁻¹).

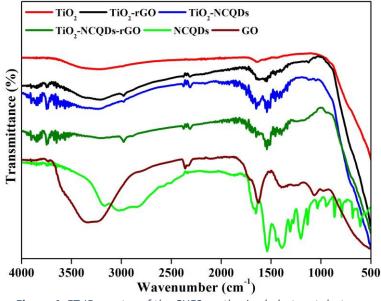


Figure 4: FT-IR spectra of the CHFS synthesised photocatalysts.

Additional peaks may be assigned for C=O/COO- (1641 cm⁻¹), C=N (1532 cm⁻¹) and C-N (1460 cm⁻¹) in the TiO₂-NCQDs supporting XPS data stated earlier. For the TiO₂-rGO related materials, the presence of oxygen functionalities on the rGO surface was supported by stretches for C=O/COO- (1627 cm⁻¹), C-OH (1396 cm⁻¹) and C-O-C vibrations (1068 cm⁻¹), which are diminished when compared to graphene oxide precursor indicating the reduced form being present. The absorption band at *ca*. 1543 cm⁻¹ may be assigned for graphene-based hybrids that relate to their skeletal vibration of the sheets [39], however, this peak may overlap the characteristic stretch for C=N of NCQDs in TiO₂-NCQDs-rGO. The broadening of peaks observed for TiO₂ nanocomposites below 1000 cm⁻¹ may correspond to TiO₂ coordinating through a combination of interactions, Ti-O-Ti and Ti-O-C, with the carbonaceous materials [40].

A photocatalyst with a bandgap that corresponds to solar light and can efficiently maximise energy utilisation would as such be highly desirable. In this regard, the light absorption properties of the as-synthesized

photocatalysts were investigated using diffuse reflectance UV-Vis spectrophotometry as profiled in **Fig. 5**, where bandgap edge values were quantified based on collected measurements.

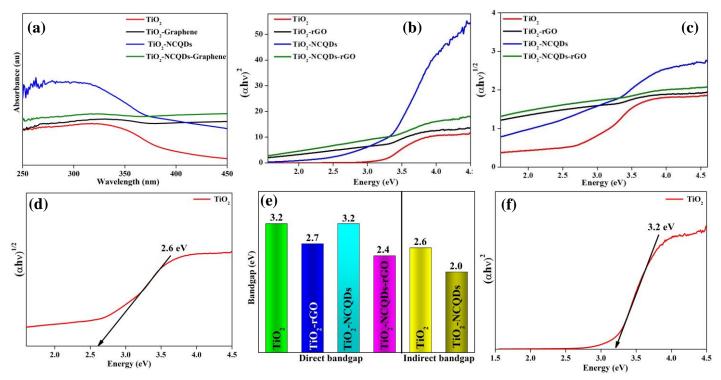


Figure 5: Optical properties of the CHFS synthesised materials: (a) UV-Vis absorption profiles, and Tauc plot profiles (α hv)ⁿ versus photon energy hv are shown for (b) direct (n=2) and (c) indirect (n =1/2) bandgap transitions of all photocatalysts materials. (e) (d and f) The resulting extrapolated band gap edge values E_g (eV) values shown for all potential photocatalysts materials. Methodology of bandgap determination for the representative sample of TiO₂ and its quantified E_g values for (d) indirect an (f) direct transition.

Characteristically, the TiO₂ (Fig. 5a) nanoparticles demonstrated a strong UV-absorption below 400 nm, however, the absorption edge shifted from the UV to visible when TiO₂ was combined with reduced graphene oxide (rGO). This phenomenon was also observed for TiO₂-NCQDs, but at a lower intensity. Significantly, it is by the incorporation of NCQDs in to TiO₂-rGO, that an appreciable increase in UV and visible radiation absorption was realised, offering the possibility of improved light-harvesting performance. The quantified bandgap edge values (Fig. 5b-c) were calculated using the Tauc plot, either for direct and indirect transitions (relationship of $[\alpha hv]^n$ versus photon energy, where n =2 for direct transition or n=1/2 for an indirect one) for each corresponding photocatalyst. The bandgap for anatase TiO₂ was calculated to be 3.2 eV, whereas the bandgap values of TiO₂-carbonaceous hybrids were noted to be substantially reduced. Extending photocatalyst absorption into the visible region of the electromagnetic spectrum by primarily narrowing the bandgap is generally regarded as one of the most difficult aspects of designing photocatalytic material. These optical studies have demonstrated that using CHFS to manufacture TiO₂ and its carbonaceous nanocomposites provided materials with improved properties when compared to anatase, such as enhanced ability to absorb visible radiation. Nevertheless, the photo-response of these modified catalysts is highly dependent on the characteristics of the type of dopant and as well as dopant optimal concentration. Crucially, fast interfacial electron transport that minimise electron-hole pair recombination (produced when a photon is absorbed), as well as the lifetime of charge separation, play key roles in the photocatalytic activity of a material.

Transient photoluminescence spectroscopy was used to assess the emission lifetime of TiO_2 and its carbonaceous nanocomposites, and to determine the charge transfer rate. To quantify the interfacial charge transfer process between TiO_2 and nitrogen-doped carbon quantum dots (NCQDs) and/or reduced graphene oxide, studies were performed using time-correlated single-photon counting (TCSPC) at an excitation wavelength of 375 nm. The normalized PL decays (**Fig. 6a**) show a decrease in the rate of decay for the nanocomposites of TiO_2 with rGO and NCQDs. The slowest decay is observed in the case of TiO_2 -NCQDs-

rGO (6.13 ns). Since PL is the result of recombination of photo-generated electrons and holes; slower decay means that this recombination is suppressed in the presence of rGO and/or NCQDs due to efficient charge separation [41].

If it is assumed that the difference of the TiO₂ emission lifetime (τ) is exclusively due to interfacial charge transfer between TiO₂ and NCQDs and/or rGO, then the rate constant for electron transfer, k_{ET}, can be determined. Previous studies[41] have estimated k_{ET} using:

$$k_{ET} = \frac{1}{\tau(TiO_2)} - \frac{1}{\tau(TiO_2 \text{ with NCQD and/or } rGO)}$$
(Equation 1)

To determine the emission lifetime, the experimental data was fitted with multiexponential decays and the average lifetimes were calculated for all samples and were used to evaluate the electron transfer rate constants (**Fig. 6b**) using **Equation 1**. Comparison of the interfacial charge transfer rate constants with trends for photocatalytic activities of the nanocomposites (**Fig. 7**) were found to be complimentary. The greatest charge transfer rates were obtained when TiO₂ was combined with NCQDs and graphene, in particular for TiO₂-NCQDs-rGO (2.6 x 10^8 s⁻¹). This reflects the photocatalytic performances of TiO₂ and its carbonaceous nanocomposites, and the influence of intimate contact between TiO₂ with NCQDs, and/or rGO.

The presence of NCQDs not only improves the composite's interfacial charge transfer but it can influence the photolytic efficiency. This was highlighted by the longer lifetimes and charge transfer rates observed for TiO₂-NCQDs and TiO₂-NCQDs-rGO when they are compared with TiO₂ and TiO₂-rGO samples.

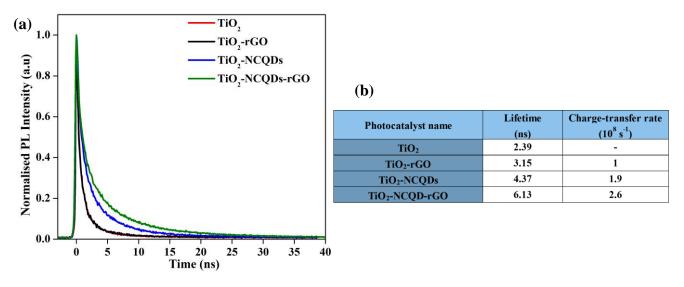


Figure 6: (a) Transient photoluminescence spectrum of TiO_2 and TiO_2 -carbonaceous nanocomposites and (b) tabulated data showing lifetime (ns) and charge transfer rate (sec⁻¹) values.

3.2 Photocatalytic activity measurements

The TiO₂ and its carbonaceous nanocomposites were examined for their ability to photodegrade a model dye pollutant (methylene blue, MB) under a solar light simulator (**Fig. 7**). The photo-response was assessed by measuring the concentration of MB at 30 minutes intervals over a total irradiation time of 180 minutes. The conversion of MB (C = final concentration, C₀ = initial concentration, (C/C₀ x100)) for each of the photocatalytic nanomaterials during total solar light exposure period was recorded (**Fig. 7a**). The photocatalytic reaction rate was determined using first-order kinetics (ln(C/C₀) = -k't) [2], where ln(C/C₀) *vs.* time exhibited a linear relationship in all the cases, and the rate constant, k', was determined from the slope. A control experiment involved the direct photodegradation of MB dye in the absence of a catalyst under the same solar irradiation conditions. Different light sources including full-spectrum solar light (> λ_{200nm}) and visible light only (> λ_{400nm}), were investigated, and the irradiation spectra are shown in **Fig. S1**.

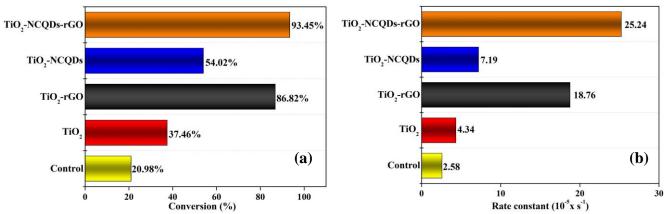


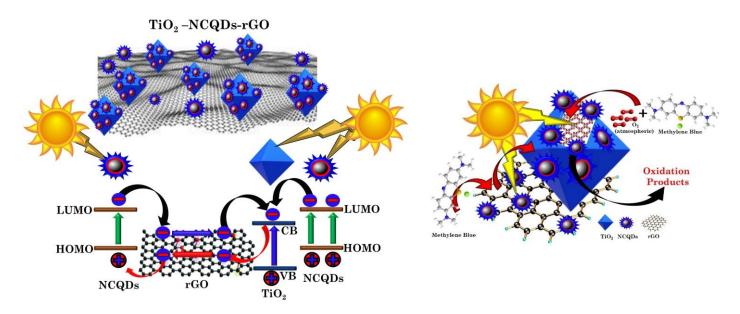
Figure 7: (a) Photo-response activities of TiO_2 and hybrid nanomaterials showing photodecomposition of methylene blue (a) the final % conversation (C= final concentration to C_0 = initial concentration, C/C₀x100) for each photocatalyst during the total solar light exposure period and, (b) first order rate constants. In both graphs, control = methylene blue.

Most of the carbonaceous TiO₂ nanocomposites (N-doped carbon quantum dots and/or reduced graphene oxide) revealed remarkable improvement on photocatalytic activity, in particular, TiO₂-NCQDs-rGO with 93% conversion of MB in the allotted 180 minutes, and significantly, a 6-fold photocatalytic rate enhancement (25.24 x10⁻⁵ s⁻¹) over that of TiO₂ alone (**Fig. 7b**). The photocatalytic activities are comparable to the lifetimes of photogenerated charge carriers (as determined from TPS and shown in **Fig. 6**), with longer lifetimes correlating with improved photocatalytic activity. Furthermore, the addition of rGO enhances the adsorption capacity towards methylene blue via either or combination of their respective anionic-cationic interactions and possible aromatic π -stacking and/or C-H- π interactions with the dye. As a consequence, the adsorbed reaction substrates can effectively react with photogenerated key active species (e.g., electrons/ holes, hydroxyl radicals, and superoxide radicals) on the surface of nanohybrid, enhancing the photoredox activity. Indeed, a TiO₂-NCQDs-rGO which exhibited highest photocatalytic activity, also had a high adsorption coefficient (**Table S1**). In all cases, the formulations of hybrids, through CHFS synthesis have proven to be more efficient in delivering enhancement of catalytic activities when compared to TiO₂ alone.

Furthermore, under visible light (λ >400 nm), the hybrid materials perform outstandingly, and the photoresponse activities (reaction rate and MB degradation) are shown in **Fig.S3**. The best performing catalyst, TiO₂-NCQDs-rGO, was selected to investigate the effect of changing light intensity (0.25, 0.5, and 1 sun) on the degradation of MB dye, and the results are shown in **Fig.S4**. It should be noted that increases in light intensity (from 0.5 to 1 sun) did not result in appreciable improvements in degradation efficiency. This indicates that at 50 mW/cm², the optimum number of photons required for efficient photocatalytic degradation was obtained (0.5 sun). Tables S2 and S3 shows the corresponding AQE values for different light sources and light intensity variations ((λ >400 nm), respectively. It was observed that an increase in light intensity corresponds to an increase in AQE, which is consistent with reported research (see supplementary information).

Tailoring optical properties of titania and controlling their electronic structure has been the subject of research primarily because it relates to optimisation of its functionalities. Hybridisation with other nanostructures such as carbon quantum dots, results in the formation of new electronic structure (type II heterojunctions) and an increase in the photo-response. This form of arrangement has been shown to not only increase the visible light harvesting potential whilst also extending the lifetime of photogenerated carriers (electrons and holes) through internal charge transfer. These phenomena are further enhanced by the substitutional doping and the addition of a heteroatom, N [42], present in all CHFS synthesised materials discussed. Graphene can also act as a co-catalyst when hybridised with titanium dioxide. With an electron conductive nanosheet and a lower Fermi stage, 2D graphene can act as an electron reservoir in accepting, transporting and shuttling the electrons produced by photoexcitation of TiO₂, enhancing charge carrier separation and transfer to participate in photo-redox processes (decomposition of methylene blue dye). Hence, multilevel/route electron transfer is possible with the hybrid structures of TiO₂ and carbonaceous materials reported herein (**Schematic 2**).

Schematic 2: Showing the mechanism of photocatalytic reaction between methylene blue and TiO₂-NCQDs-rGO hybrid. A type II heterojunction is formed when TiO₂ and NCQDs combine. When excited by full solar light irradiation, electrons (e⁻) move from the highest-occupied molecular orbital (HOMO) of NCQDs and the valence band (VB) of TiO₂ to the lowest-occupied molecular orbital (LUMO)/CB. Literature suggests that anatase prefers h⁺ mediated photocatalysis for dye photodegradation. As a result, the photoinduced holes (h⁺) can oxidise water molecules, producing hydroxyl free radicals (•OH), which are very reactive in oxidising of MB. The excited electrons migrate into the π -orbital of the rGO layer, which serves as a conducting 2D network, where excited electrons can freely move to combine with the dissolved oxygen, generating oxygen radicals (•O₂⁻) that can also react with MB. The movement of photoexcited electrons from the LUMO level of NCQDs to the rGO layer is also expected. Experiments with various light sources show that TiO₂ absorbs UV and visible light (due to N-doping creating allowed electronic states between VB and CB of titania). Overall, the material combination enables efficient charge separation and rapid charge carrier diffusion, resulting in a low charge recombination rate.



3.3 Density functional theory (DFT) calculations

To gain further insight into how the rGO layer (or carbon quantum dots) affects the position of the electronic energy band edges of the anatase system, and what charge transfer occurs if any, an idealised TiO_2 -graphene model was subjected to density functional theory (DFT) calculations using a plane-wave basis set and a hybrid functional for exchange and correlation. A trade-off between model size and accuracy of the calculation of the electronic properties was necessary, thus, utilisation of simple slab models to simulate the interface between anatase and graphene, was chosen. To carry out the hybrid DFT calculations tractably, the size of the model was restricted to 100-120 atoms to allow fitting of the graphene layers on the slabs with periodic boundary conditions and allowing large strains (greater than 5 %) being applied to the graphene layer.

This model of highly strained graphene layers adhered to effectively an infinite in extent 2D slabs of TiO₂, is a limited simulation of the as-synthesised nanocomposites. Nevertheless, substantial insight can still be gained into the expected charge transfer and modulation of the energy band edges of anatase by the rGO and/or carbon dots that allowed for enhancement of photocatalytic properties [43]. Realistically, larger models could be designed using, for example, interatomic force fields and/or semiclassical techniques [50, 51], however, such methods tend to include empirical parameters which have their own disadvantages. As such, the approach chosen herein provides a significantly more accurate account of the electronic energy levels, where analysis of the variation in said energy levels when graphene is present allows for identification of any possible improvements in photocatalytic properties.^{*}

The idealised model, chosen as the more energetically favourable for anatase, consisted of slabs of anatase TiO_2 cut along the low-index (001), (101) and (100) surfaces [46]. The slabs were relaxed fully, and the electron ionisation potentials, *I*, computed. Knowledge of *I*, the energy required to excite an electron from the

^{*}A more complete study, using larger slab models is underway and will be reported in the future.

valence band within the slab to the vacuum level outside, allows for referencing the valence band edges to the vacuum level, thus placing the band on an absolute scale [26,47]. Combined with the band gap of bulk anatase, allows for determination of where the valence band maximum (VBM) and conduction band minimum (CBM) lie relative to particular redox potentials, for example, water oxidation and reduction potentials for photocatalytic water splitting. Once *I* has been computed, a layer of graphene was added to the TiO₂ slab, and the combined TiO₂-graphene system was allowed to relax. Then, *I* for the combined system was calculated to determine the effect of graphene on the band edges. Moreover, by analysing the partial charge densities and spin density, the charge transfer between graphene and TiO₂ was shown to occur.

As shown below, a small negative charge density is transferred from the TiO_2 slab to the graphene layer, meaning that the VBM states are confined within the graphene layer and the CBM states within the TiO_2 layer. Furthermore, *I* increases substantially within the graphene layers, so that the VBM is pushed downwards (to more negative energies) on the absolute scale, whilst *I* within the interior of the TiO_2 slab is decreased, to facilitate CBM being pushed upwards on the absolute scale. Crucially, both of these effects, i.e. the separation of charge and the pushing down (up) of the VBM (CBM) on the absolute energy scale, will be beneficial for photocatalytic activity [48]. The charge separation represses unwanted charge recombination, while the energy levels shift to more favourable positions relative to the water redox potentials. These two effects constitute the main results for the computational study here within. Although the slab model is a simplification with respect to the systems for which the measurements were performed, nevertheless it was expected that both effects should occur in our samples as a significant source for improved photocatalytic activity of the composite systems. We will now present the detailed results of the DFT calculations.

After relaxing the atomic coordinates, the lattice parameters, a = 3.789 Å and c = 9.613 Å, were determined within ~1% of the experimental values for the tetragonal anatase structure [49]. The calculated band gap of anatase was 3.40 eV and in good agreement with previous theoretical and experimental studies [47,50]. The relaxed C-C bond length in graphene was found to be 1.415 Å, in excellent agreement with the experimental C-C bond in graphite and graphene (1.42 Å). The (001), (100) and (101) surfaces of anatase TiO₂ were relaxed for the slab model approach (Figure). As the surfaces are non-polar, minor relaxations were observed on the surface. The calculated surface energies: (001) 1.21 Jm⁻², (100) 0.73 Jm⁻² and (101) 0.16 Jm⁻², which follow the observed trend were in reasonable agreement with standard DFT [51,52].

Shown in the lower panes of **Figure**, are the relaxed slabs with the added graphene layers. After relaxation, the graphene layers were found to have rumpled very slightly, with movements of the C atoms of ~0.01-0.02 Å. The graphene layers relaxed to distances of 3.76 Å, 3.51 Å and 3.52 Å from the (001), (100) and (101) anatase surfaces, respectively. The potential energy surface proved to be quite shallow; consequently, the relatively large convergence criterion in interatomic forces of 0.025 eVÅ⁻¹ was chosen for these relaxations. Interestingly, for the (100) surface, a second local minimum in the potential energy surface was found that had a total energy slightly lower (by about 5 meV per atom) than the relaxed surface shown in **Figure**. This result showed larger rumpling by an order of magnitude in the graphene layer.

The calculated planar electrostatic potentials averaged over the cells in the direction perpendicular to the slab surfaces are shown in **Figure**, along with the computed values of ΔV (indicated by arrows). From the relaxation of bulk anatase, it was determined that $E_{VBM} = -0.45$ eV and $E_g = 3.40$ eV. It was therefore found for the cases without graphene, the surface cuts (001), (100) and (101) have I = 7.2 eV, 7.3 eV and 7.5 eV respectively, and A = -3.8 eV, -3.9 eV and -4.1 eV respectively.

The change in the potential when graphene was added was quite pronounced. Two clear potential wells were observed, one deeper well close to the surface of the slab, and a broad, shallower well in the interior of the slab, relative to the surface. As the offset in potential between the slab interior and the vacuum can be related to the position of the VBM on an absolute scale, we see that the energy bands are pushed down at the surface, and up in the interior of the slabs. For the (001) surface, the shift of the bands to more negative energy at the surface when graphene is added is -0.1 eV, and the shift to more positive energies in the interior is 0.3 eV. For the (100) surface, the shifts are -0.7 eV and +0.3 eV, and for the (101) surface they are -0.8 eV and +0.8 eV. For the (001) surface, the deeper well at the surface when graphene is added has a shallow part, then a sharp

deepening, in contrast to the two other cases studied. It was observed that in the case of the (001) surface, the TiO_2 surface below the anatase layer relaxed more substantially than in the other cases studied. *

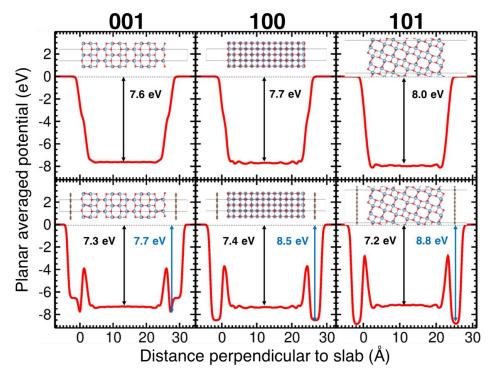


Figure 8: Calculated planar averaged potential along the direction perpendicular to the slabs, given relative to the value in the vacuum gap. The distance perpendicular to the slabs is given relative to one of the edges of the pure TiO_2 slab for each of the facets modelled, which are (left to right): (001), (100) and (101). In the upper panes, the results are given for the pure TiO_2 slabs, while the lower panes show the results for the corresponding slabs with single layers of graphene added above and below. The arrows indicate the computed offsets in potential between the vacuum layer and slab interiors (ΔV). The slab models are shown in each pane, aligned with their planar electrostatic potential. Ti atoms are represented by blue spheres, oxygen atoms by red spheres and carbon atoms by brown spheres.

A clear shift downwards in energy of the bands at the graphene layer was observed, and upwards in the TiO₂ layer, for the combined systems. Shown in **Fig. 9** are the spin density (**Figure b**) and the partial charge densities for the states at the VBM (**Figure c**) and CBM (**Figure d**), for the slab containing the (100) anatase surface (for the other cases, similar results were observed). A charge transfer from the TiO₂ slab to the graphene layer of ~0.1 electrons (in total in the slab model) was also observed. Moreover, as seen from the partial charge densities associated with the VBM and CBM, the states at the VBM are predominately concentrated in the graphene layer, while the states at the CBM are localised in the interior of the TiO₂ slab, on the Ti atoms with charge densities similar to the shapes of d orbitals. It can therefore be concluded that the effective VBM of the combined system is pushed downwards on an absolute energy scale, whilst the CBM is pushed upwards. This shift in band edges results in a more favourable alignment with the water redox potentials, thus favouring improved catalytic activity [26,47]. The charge separation will also favour photocatalytic processes, as recombination will be suppressed. The combination of these two effects is significant and should see considerable improvement in photocatalytic performance, due to the presence of the graphene on the TiO₂ slab.

Whilst this model of a TiO_2 -graphene nanocomposite, as discussed earlier, is somewhat crude, nevertheless, these electronic effects, calculated with a sophisticated first principles method, can be expected in real samples. The absolute calculated shifts are not expected to be accurate when compared with experimental results but are indicative of trends that are crucial for improved catalytic performance. The modelling has demonstrated that the combination of anatase and graphene results in energy band shifts and charge separation that significantly benefits photocatalytic performance, and it can be expected that these effects will also be present to some degree when other carbon structures are used, such as quantum dots and rGO layers.

^{*} We are investigating these relaxation effects further, but a full analysis is beyond the scope of the current study.

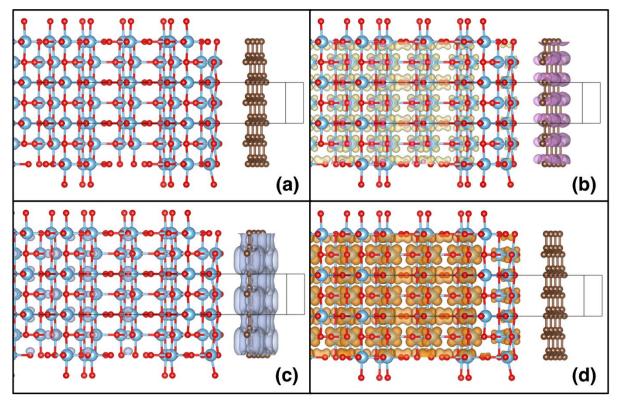


Figure 9: The spin and charge densities for the relaxed combined anatase (100) surface and graphene layer slab. (a) The relaxed (100) anatase surface with the graphene layer above. The direction perpendicular to the slab is to the right. (b) The calculated spin density (purple and yellow isosurface with a density of $10^{-4} e/Å^3$). (c) The partial charge density associated with the states at the top of the valence band (grey isosurface with a density of $10^{-3} e/Å^3$). (d) the partial charge density associated with the state at the bottom of the conduction band (orange isosurface with a density of $10^{-3} e/Å^3$). The the state by blue spheres, oxygen atoms by red spheres and carbon atoms by brown spheres.

4. Conclusions

Continuous hydrothermal flow synthesis enabled the rapid and efficient production of homogeneous nano-TiO₂ hybrid composites with nitrogen-doped carbon quantum dots and/or reduced graphene oxide with outstanding photocatalytic activities. Titania particles exhibited a particle size of < 5 nm, high crystallinity (anatase form) and large surface area (typically > $232 \text{ m}^2 \text{ g}^{-1}$), NCQDs exhibited spherical morphology (mean particle size ca. 3.4 nm) with a graphitic core arrangement, and the graphene (rGO) sheets consisted of 5-8 layers. The highlight of the TiO₂ hybrid nanocomposites remarkable photocatalytic activity was TiO₂-NCQDs-rGO providing 93% conversion of methylene blue with a 6-fold photocatalytic rate enhancement $(25.24 \text{ x}10^{-5} \text{ s}^{-1})$ over that of TiO₂ alone under solar light simulation, was facilitated by longer emission lifetimes (6.13 ns), faster charge transfer rates ($2.6 \times 10^8 \text{ s}^{-1}$), and a narrowed band gap. Changes in the bandgap are supported by Hybrid DFT calculations based on a model system that indicated a shift in band edges to a more favourable alignment between anatase with graphene, and recombination would be suppressed. The combination of these two effects account for the considerable improvement in photocatalytic performance identified for TiO₂-rGO example (emission lifetime, 3.15 ns), charge transfer rate (1.0×10^8) s^{-1}). The model proposed for TiO₂-NCQDs-rGO would also incorporate this, facilitated by NCQDs further increasing the photosensitivity, the emission lifetimes and charge transfer rates via multilevel electron transfer thus significantly enhancing photocatalytic activity.

Effective nano level engineering coupled with a successful synthetic technique that delivers close interfacial interaction and tailored surface chemistry are thus critical for achieving the optimum performance and efficiency in terms of the charge carrier transfer to the surface of the photocatalyst facilitating relevant photoredox reactions to take place. These hybrid structures reported here have potential to be explored in alternative fields including energy storage, hydrogen generation, or environmental applications.

Conflict of interest

The authors have no conflicts of interest to declare.

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Author contributions:

IAB: Investigation, writing original draft; JB: Theoretical calculations and associated methodology, writing original draft; KGN: photocatalytic experiments under LED light, TH: TEM analysis; TS, ST and AR: Transient photoluminescence spectroscopy analysis; DJM: XPS analysis; NPP: XRD characterisation data; SAN and MT: UV-VIS spectrophotometry absorption measurements of the photocatalysts powders; SD and SK: conceptualisation, facilities and resources; SK; project management/supervision, methodology, writing original draft; All authors reviewed and edited the paper.

References:

- [1] U. Nations, The United Nations World Water Development Report 2019: Leaving no one behind, 2019.
- [2] J.B.M. Goodall, S. Kellici, D. Illsley, R. Lines, J.C. Knowles, J.A. Darr, Optical and photocatalytic behaviours of nanoparticles in the Ti-Zn-O binary system, RSC Adv. 4 (2014) 31799–31809. https://doi.org/10.1039/c3ra48030e
- [3] P. Kar, S. Zeng, Y. Zhang, E. Vahidzadeh, A. Manuel, R. Kisslinger, K.M. Alam, U.K. Thakur, N. Mahdi, P. Kumar, K. Shankar, High rate CO₂ photoreduction using flame annealed TiO₂ nanotubes, Appl. Catal. B Environ. 243 (2019) 522–536. https://doi.org/10.1016/j.apcatb.2018.08.002.
- [4] [a] S. Chen, T. Takata, K. Domen, Particulate photocatalysts for overall water splitting, Nat. Rev. Mater. 2 (2017) 1–17. https://doi.org/10.1038/natrevmats.2017.50. [b] R. Su, Designer titaniasupported Au–Pd nanoparticles for efficient photocatalytic hydrogen productio, ACS Nano, 8(4) (2014) 3490–3497. https://doi.org/10.1021/nn500963m.
- [5] C.F. Goodeve, J.A. Kitchener, Photosensitisation by titanium dioxide, Trans. Faraday Soc. 34 (1938) 570–579. https://doi.org/10.1039/TF9383400570.
- [6] Y. Zhang, K. Hawboldt, L. Zhang, J. Lu, L. Chang, A. Dwyer, Carbonaceous nanomaterial-TiO₂ heterojunctions for visible-light-driven photocatalytic degradation of aqueous organic pollutants, Appl. Catal. A Gen. 630 (2022) 118460. https://doi.org/10.1016/j.apcata.2021.118460.
- [7] E. Kusiak-Nejman, A.W. Morawski, TiO₂/graphene-based nanocomposites for water treatment: A brief overview of charge carrier transfer, antimicrobial and photocatalytic performance, Appl. Catal. B Environ. 253 (2019) 179–186. https://doi.org/10.1016/j.apcatb.2019.04.055.
- [8] B. Yao, H. Huang, Y. Liu, Z. Kang, Carbon Dots: A Small Conundrum, Trends Chem. 1 (2019) 235– 246. https://doi.org/10.1016/j.trechm.2019.02.003.
- [9] H. Luo, S. Dimitrov, M. Daboczi, J.S. Kim, Q. Guo, Y. Fang, M.A. Stoeckel, P. Samorì, O. Fenwick, A.B. Jorge Sobrido, X. Wang, M.M. Titirici, Nitrogen-doped carbon dots/TiO₂ nanoparticle composites for photoelectrochemical water oxidation, ACS Appl. Nano Mater. 3 (2020) 3371–3381. https://doi.org/10.1021/acsanm.9b02412.
- [10] N. Zhang, M.Q. Yang, S. Liu, Y. Sun, Y.J. Xu, Waltzing with the versatile platform of graphene to synthesize composite photocatalysts, Chem. Rev. 115 (2015) 10307–10377. https://doi.org/10.1021/acs.chemrev.5b00267.
- [11] K.R. Reddy, M. Hassan, V.G. Gomes, Hybrid nanostructures based on titanium dioxide for enhanced photocatalysis, Appl. Catal. A Gen. 489 (2015) 1–16. https://doi.org/10.1016/j.apcata.2014.10.001.
- [12] X. Yang, D. Wang, Photocatalysis: From fundamental principles to materials and applications, ACS Appl. Energy Mater. 1 (2018) 6657–6693. https://doi.org/10.1021/acsaem.8b01345.
- [13] J.A. Darr, J. Zhang, N.M. Makwana, X. Weng, Continuous hydrothermal synthesis of inorganic

nanoparticles: applications and future directions, Chem. Rev. 117 (2017) 11125–11238. https://doi.org/10.1021/acs.chemrev.6b00417.

- [14] V. Middelkoop, C.J. Tighe, S. Kellici, R.I. Gruar, J.M. Perkins, S.D.M. Jacques, P. Barnes, J.A. Darr, Imaging the continuous hydrothermal flow synthesis of nanoparticulate CeO₂ at different supercritical water temperatures using in situ angle-dispersive diffraction, J. Supercrit. Fluids. 87 (2014) 118–128. https://doi.org/10.1016/j.supflu.2013.12.022.
- [15] S. Kellici, J. Acord, A. Vaughn, N.P. Power, D.J. Morgan, T. Heil, S.P. Facq, G.I. Lampronti, Calixarene assisted rapid synthesis of silver-graphene nanocomposites with enhanced antibacterial activity, ACS Appl. Mater. Interfaces. 8 (2016) 19038–19046. https://doi.org/10.1021/acsami.6b06052.
- [16] U. Alli, S.J. Hettiarachchi, S. Kellici, Chemical functionalisation of 2D materials by batch and continuous hydrothermal flow synthesis, Chem. - A Eur. J. 26 (2020) 6447–6460. https://doi.org/10.1002/chem.202000383.
- [17] S. Kellici, J. Acord, J. Ball, H.S. Reehal, D. Morgan, B. Saha, A single rapid route for the synthesis of reduced graphene oxide with antibacterial activities, RSC Adv. 4 (2014) 14858–14861. https://doi.org/10.1039/c3ra47573e.
- [18] I.A. Baragau, N.P. Power, D.J. Morgan, R.A. Lobo, C.S. Roberts, M.M. Titirici, V. Middelkoop, A. Diaz, S. Dunn, S. Kellici, Efficient continuous hydrothermal flow synthesis of carbon quantum dots from a targeted biomass precursor for on-off metal ions nanosensing, ACS Sustain. Chem. Eng. 9 (2021) 2559–2569. https://doi.org/10.1021/acssuschemeng.0c08594.
- [19] I.A. Baragau, N.P. Power, D.J. Morgan, T. Heil, R.A. Lobo, C.S. Roberts, M.M. Titirici, S. Dunn, S. Kellici, Continuous hydrothermal flow synthesis of blue-luminescent, excitation-independent nitrogen-doped carbon quantum dots as nanosensors, J. Mater. Chem. A. 8 (2020) 3270–3279. https://doi.org/10.1039/c9ta11781d.
- [20] S. Kellici, J. Acord, N.P. Power, D.J. Morgan, P. Coppo, T. Heil, B. Saha, Rapid synthesis of graphene quantum dots using a continuous hydrothermal flow synthesis approach, RSC Adv. 7 (2017) 14716–14720. https://doi.org/10.1039/C7RA00127D.
- [21] S. Kellici, J. Acord, K.E. Moore, N.P. Power, V. Middelkoop, D.J. Morgan, T. Heil, P. Coppo, I.A. Baragau, C.L. Raston, Continuous hydrothermal flow synthesis of graphene quantum dots, React. Chem. Eng. 3 (2018) 949–958. https://doi.org/10.1039/c8re00158h.
- [22] G. Kresse, J. Hafner, Ab initio molecular dynamics for liquid metals, Phys. Rev. B. 47 (1993) 558– 561. https://doi.org/10.1103/PhysRevB.47.558.
- [23] G. Kresse, J. Furthmüller, Efficiency of ab-initio total energy calculations for metals and semiconductors using a plane-wave basis set, Comput. Mater. Sci. 6 (1996) 15–50. https://doi.org/10.1016/0927-0256(96)00008-0.
- [24] G. Kresse, J. Hafner, Ab initio molecular-dynamics simulation of the liquid-metalamorphoussemiconductor transition in germanium, Phys. Rev. B. 49 (1994) 14251–14269. https://doi.org/10.1103/PhysRevB.49.14251.
- [25] A. V. Krukau, O.A. Vydrov, A.F. Izmaylov, G.E. Scuseria, Influence of the exchange screening parameter on the performance of screened hybrid functionals, J. Chem. Phys. 125 (2006). https://doi.org/10.1063/1.2404663.
- [26] D.O. Scanlon, C.W. Dunnill, J. Buckeridge, S.A. Shevlin, A.J. Logsdail, S.M. Woodley, C.R.A. Catlow, M.J. Powell, R.G. Palgrave, I.P. Parkin, G.W. Watson, T.W. Keal, P. Sherwood, A. Walsh, A.A. Sokol, Band alignment of rutile and anatase TiO₂, Nat. Mater. 12 (2013) 798–801. https://doi.org/10.1038/nmat3697.
- [27] P.E. Blöchl, Projector augmented-wave method, Phys. Rev. B. 50 (1994) 17953–17979. https://doi.org/10.1103/PhysRevB.50.17953.
- [28] Hendrik J. Monkhorst and James D. Pack, Special points for Brillonin-zone integrations*, Phys. Rev. B. 13 (1976) 5188–5192. https://doi.org/https://doi.org/10.1039/C3RA48030E.
- [29] S. Project, SAINT Surface and Interface Toolkit: For the Materials and Molecular Modelling Community, Https://Saint.Chem.Ucl.Ac.Uk/. (n.d.).
- [30] M. Peressi, N. Binggeli, A. Baldereschi, Band engineering at interfaces: Theory and numerical experiments, J. Phys. D. Appl. Phys. 31 (1998) 1273–1299. https://doi.org/10.1088/0022-3727/31/11/002.

- [31] A. Walsh, K.T. Butler, Prediction of electron energies in metal oxides, Acc. Chem. Res. 47 (2014) 364–372. https://doi.org/10.1021/ar400115x.
- [32] K.T. Butler, C.H. Hendon, A. Walsh, Electronic chemical potentials of porous metal-organic frameworks, J. Am. Chem. Soc. 136 (2014) 2703–2706. https://doi.org/10.1021/ja4110073.
- [33] N.R. Khalid, A. Majid, M.B. Tahir, N.A. Niaz, S. Khalid, Carbonaceous-TiO₂ nanomaterials for photocatalytic degradation of pollutants: A review, Ceram. Int. 43 (2017) 14552–14571. https://doi.org/10.1016/j.ceramint.2017.08.143.
- [34] K. Thompson, J. Goodall, S. Kellici, J.A. Mattinson, T.A. Egerton, I. Rehman, J.A. Darr, Screening tests for the evaluation of nanoparticle titania photocatalysts, J. Chem. Technol. Biotechnol. 84 (2009) 1717–1725. https://doi.org/10.1002/jctb.2237.
- [35] J.S. Lee, K.H. You, C.B. Park, Highly photoactive, low bandgap TiO₂ nanoparticles wrapped by graphene, Adv. Mater. 24 (2012) 1084–1088. https://doi.org/10.1002/adma.201104110.
- [36] N.S. Allen, N. Mahdjoub, V. Vishnyakov, P.J. Kelly, R.J. Kriek, The effect of crystalline phase (anatase, brookite and rutile) and size on the photocatalytic activity of calcined polymorphic titanium dioxide (TiO₂), Polym. Degrad. Stab. 150 (2018) 31–36. https://doi.org/10.1016/j.polymdegradstab.2018.02.008.
- [37] W. Vallejo, A. Rueda, C. Díaz-Uribe, C. Grande, P. Quintana, Photocatalytic activity of graphene oxide-TiO₂ thin films sensitized by natural dyes extracted from Bactris guineensis, R. Soc. Open Sci. 6 (2019). https://doi.org/10.1098/rsos.181824.
- [38] and C.P. Alexander Naumkin, Anna Kraut-Vass, Stephen Gaarenstroom, NIST X-ray Photoelectron Spectroscopy Database, NIST Stand. Ref. Database 20, Version 4.1. (2012).
- [39] H. Zhang, X. Lv, Y. Li, Y. Wang, J. Li, P25-graphene composite as a high performance photocatalyst, ACS Nano. 4 (2010) 380–386. https://doi.org/10.1021/nn901221k.
- [40] Y. Zhou, S. Yang, D. Fan, J. Reilly, H. Zhang, W. Yao, J. Huang, Carbon quantum dot/TiO₂ nanohybrids: efficient photocatalysts for hydrogen generation via intimate contact and efficient charge separation, ACS Appl. Nano Mater. 2 (2019) 1027–1032. https://doi.org/10.1021/acsanm.8b02310.
- [41] L. Liu, Z. Liu, A. Liu, X. Gu, C. Ge, F. Gao, L. Dong, Engineering the TiO₂-graphene interface to enhance photocatalytic H₂ production, ChemSusChem. 7 (2014) 618–626. https://doi.org/10.1002/cssc.201300941.
- [42] V. Štengl, D. Popelková, P. Vláčil, TiO₂-graphene nanocomposite as high performace photocatalysts, J. Phys. Chem. C. 115 (2011) 25209–25218. https://doi.org/10.1021/jp207515z.
- [43] J. Schneider, M. Matsuoka, M. Takeuchi, J. Zhang, Y. Horiuchi, M. Anpo, D.W. Bahnemann, Understanding TiO₂ Photocatalysis Mechanisms and Materials, Chem. Rev. 114 (2014) 9919–9986.
- [44] X.J. Han, L. Bergqvist, P.H. Dederichs, H. Müller-Krumbhaar, J.K. Christie, S. Scandolo, P. Tangney, Polarizable interatomic force field for TiO₂ parametrized using density functional theory, Phys. Rev. B Condens. Matter Mater. Phys. 81 (2010) 1–9. https://doi.org/10.1103/PhysRevB.81.134108.
- [45] R.S. Kavathekar, P. Dev, N.J. English, J.M.D. MacElroy, Molecular dynamics study of water in contact with the TiO₂ rutile-110, 100, 101, 001 and anatase-101, 001 surface, Mol. Phys. 109 (2011) 1649–1656. https://doi.org/10.1080/00268976.2011.582051.
- [46] G. Liu, H.G. Yang, J. Pan, Y.Q. Yang, G.Q.M. Lu, H.M. Cheng, Titanium dioxide crystals with tailored facets, Chem. Rev. 114 (2014) 9559–9612. https://doi.org/10.1021/cr400621z.
- [47] J. Buckeridge, K.T. Butler, C.R.A. Catlow, A.J. Logsdail, D.O. Scanlon, S.A. Shevlin, S.M. Woodley, A.A. Sokol, A. Walsh, Polymorph engineering of TiO₂: Demonstrating how absolute reference potentials are determined by local coordination, Chem. Mater. 27 (2015) 3844–3851. https://doi.org/10.1021/acs.chemmater.5b00230.
- [48] P. Deák, J. Kullgren, B. Aradi, T. Frauenheim, L. Kavan, Water splitting and the band edge positions of TiO₂, Electrochim. Acta. 199 (2016) 27–34. https://doi.org/10.1016/j.electacta.2016.03.122.
- [49] D.T. Cromer, K. Herrington, The structures of anatase and rutile, J. Am. Chem. Soc. 77 (1955) 4708– 4709. https://doi.org/10.1021/ja01623a004.
- [50] B. Kraeutler, A.J. Bard, Heterogeneous photocatalytic decomposition of saturated carboxylic acids on TiO₂ powder. decarboxylative route to alkanes, J. Am. Chem. Soc. 19 (1978) 5985–5991.
- [51] M. Lazzeri, A. Vittadini, A. Selloni, Structure and energetics of stoichiometric TiO₂ anatase surfaces,

Phys. Rev. B - Condens. Matter Mater. Phys. 63 (2001) 1554091–1554099. https://doi.org/10.1103/PhysRevB.63.155409.

[52] P.R. B, M. Lazzeri, A. Vittadini, A. Selloni, Erratum: Structure and energetics of stoichiometric anatase surfaces [Phys. Rev. B 63, 155409 (2001)], Phys. Rev. B - Condens. Matter Mater. Phys. 65 (2002) https://doi.org/10.1103/PhysRevB.65.119901.