

# Origins of ductile plasticity in a polycrystalline gallium arsenide during scratching: MD simulation study

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## Abstract

This paper used molecular dynamics simulation to reveal the origins of the ductile plasticity in polycrystalline gallium arsenide (GaAs) during its nanoscratching. Velocity-controlled nanoscratching tests were performed with a diamond tool to study the friction-induced deformation behaviour of polycrystalline GaAs. Cutting temperature, sub-surface damage depth, cutting stresses, the evolution of dislocations

and the subsequent microstructural changes were extracted from the simulation. The simulated MD data indicated that the deformation of polycrystalline GaAs is accompanied by dislocation nucleation in the grain boundaries (GBs) leading to the initiation of plastic deformation. Furthermore, the  $1/2\langle 110 \rangle$  is the main type of dislocation responsible for ductile plasticity in polycrystalline GaAs. The magnitude of cutting forces and the extent of sub-surface damage were both observed to reduce with an increase in the scratch velocity whereas the cutting temperature scaled with the cutting velocity. As for the depth of the scratch, an increase in its magnitude increased the cutting forces, temperature and damage-depth. A phenomenon of fluctuation from wave crests to wave troughs in the cutting forces was observed only during the cutting of polycrystalline GaAs and not during the cutting of single-crystal GaAs.

**Keywords:** Polycrystalline gallium arsenide; MD simulation; Grain boundary; Dislocation nucleation

## **1 Introduction**

Gallium arsenide (GaAs) has emerged as a favourable choice as a III-V direct bandgap semiconductor due to its applications in 5G communication devices [1]. GaAs (which resides in a Zinc-blende structure) possesses superior properties to silicon, for instance, GaAs has a higher saturated electron velocity and higher electron mobility, allowing GaAs transistors to function swiftly at frequencies over 250 GHz. Owing to their wider energy bandgap, GaAs devices are relatively insensitive to overheating which makes them less noisy while operating at higher frequencies in electronic circuits and that is

where they outbid silicon devices [2]. GaAs can be grown as a single crystal using methods such as the vertical gradient freeze method, the Bridgman-Stockbarger technique, or the Liquid encapsulated Czochralski growth process [3][4]. Parallel to this, the films of polycrystalline GaAs can be grown by chemical vapour deposition (by annealing an amorphously grown film) [5], or by using molecular beam epitaxy (MBE) [6]. Remarkably, it was found that, comparing the application fields between the single crystal GaAs and polycrystalline GaAs, the single crystal GaAs is widely applied in the wireless communication aiming to offer the data communication between base station and users [7]. However, the polycrystalline GaAs is intensively employed in fiber optic communication aiming to complete the centralized transmission of a large number of user's data [8]. As opposed to the single crystal GaAs, solar cells of polycrystalline GaAs thin film based can much better reach the demand in the high efficiency (20% AM1.5 with average grain sizes  $< 1 \text{ mm}^2$ ) [9], thin [10], light [11] and flexibility [12]. In the field of imaging detectors [13], microwave [14] and optoelectronic devices [15], the polycrystalline GaAs is allowed wider application due to its low cost than processing cost of single crystal GaAs. For the polycrystalline GaAs based nanoscale devices in above applications, multiplex 2D or 3D free-form nanostructures are often required. Subsequently, the investigation of ductile plasticity mechanism become significant during nanocutting polycrystalline GaAs. The deformation mechanisms of polycrystal material which dominated by grain boundaries and dislocations are widely studied and reported in many previous studies[16][17][18]. However, the mechanism was studied by concentrating on the dislocations nucleation inside of the grains. The

origin of incipient dislocations site and fundamental reasons of ductile plasticity in a polycrystalline gallium arsenide during scratching is unknown. Therefore, these knowledge gaps prompt the authors to investigate the origin of ductile plasticity in a polycrystalline gallium arsenide by establishing an extreme scratching conditions spanning from 0 nm to 2 nm cut of depth via effective molecular dynamics (MD) simulation technique [19]. The remaining sections of the paper discuss the scratch forces, sub-surface damage, peak cutting temperature, cutting stresses in a polycrystalline substrate benchmarked against a single crystal GaAs substrate.

## **2 MD simulation methodology**

### *2.1 MD simulation model*

An orthogonal MD simulation nanoscratching model for cutting a polycrystalline GaAs was developed using the Voronoi site-rotation algorithm (see figure 1 and figure 2(a)). The diamond scratching tool was modelled with a negative rake angle of  $-25^\circ$  and a clearance angle of  $10^\circ$ . The polycrystalline GaAs workpiece was built by employing the AtomsK software [20] using the Voronoi algorithm [21][22]. The Voronoi site-rotation method generates a polycrystalline structure by joining the normals of the line of random discrete points at the crystal boundary and containing the growing random oriented crystal seeds. As shown in Fig. 2(b), the polycrystalline GaAs workpiece was divided into three regions of atoms i.e. Newtonian atoms (yellow and red color atoms), thermostat atoms (blue color atoms) and boundary atoms (green color atoms). Prior to cutting, the GaAs workpiece (containing 14 grains) was equilibrated for about 100 ps

using the Nose-Hoover method [23] in LAMMPS (Large-scale atomic/molecular massively parallel simulator) [24]. Visualization and analysis were performed using Visual molecular dynamics (VMD) [25] and Open Visualization Tool (OVITO) softwares [26]. The detailed parameters used for the model development are shown in Table 1.

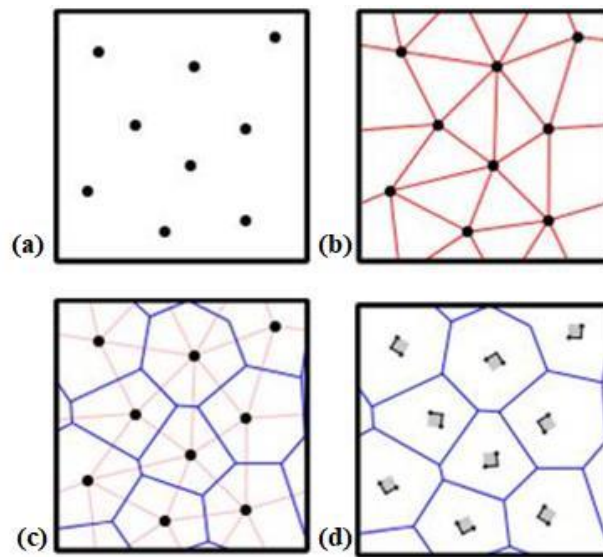
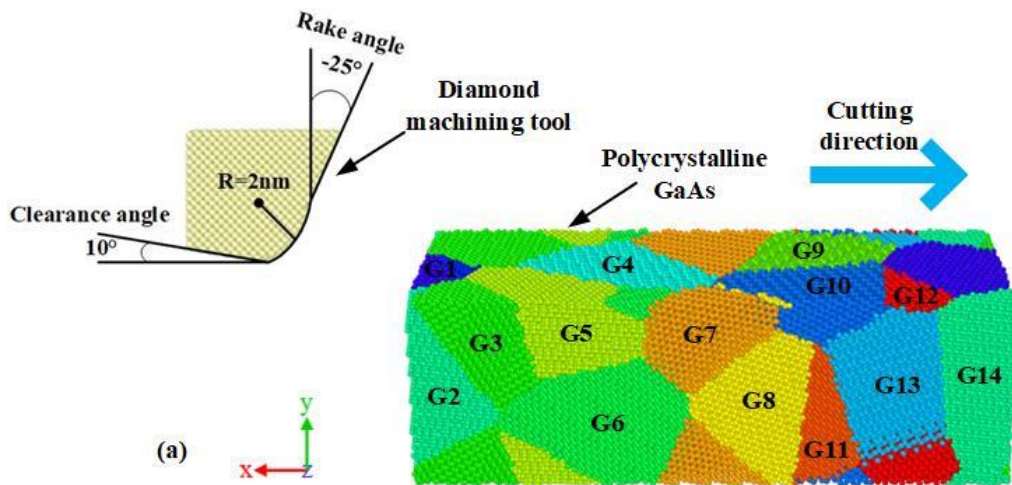
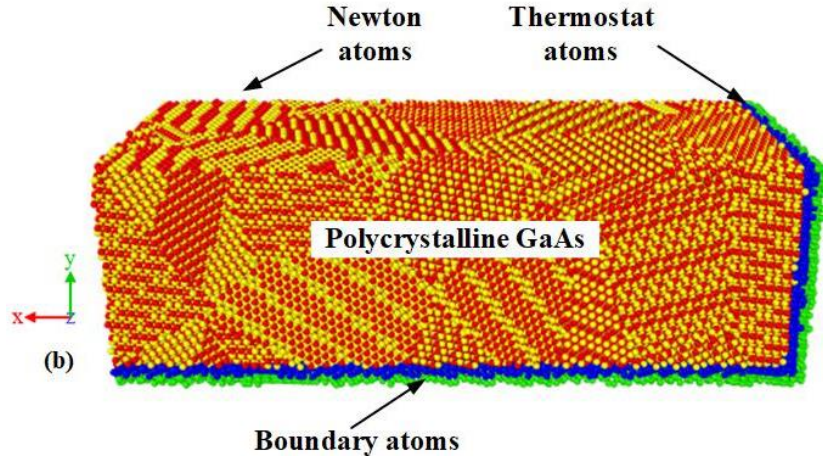


Fig. 1. Voronoi site-rotation illustration showing random points.





**Fig. 2.** Nanoscratching model of polycrystalline GaAs showing different grains and tool description.

**Table 1:** Simulation parameters used to develop the MD simulation model

GaAs substrate dimensions	30.8 nm × 10.0 nm × 13.4 nm (X, Y and Z direction)
Number of atoms in the polycrystalline GaAs	184285 (total 14 grains)
Scratching tool	Diamond cutting tool (rigid)
Number of atoms in the tool	12085
Tool rake angle	-25°
Tool clearance angle	10°
Tool edge radius	2 nm
Equilibrium lattice constant of GaAs	5.78 Å (Zinc blende lattice structure)
Diamond lattice constant	3.57 Å (Diamond cubic lattice structure)
Width of cut	2.86 nm
Depth of cut	Was varied (0.5 nm, 1 nm and 2 nm)
Scratch velocity	Was varied (100 m/s, 200 m/s and 400 m/s)
Scratching distance	12 nm

Equilibration temperature	300 K
Timestep	1 fs

## *2.2 Selection of potential energy function*

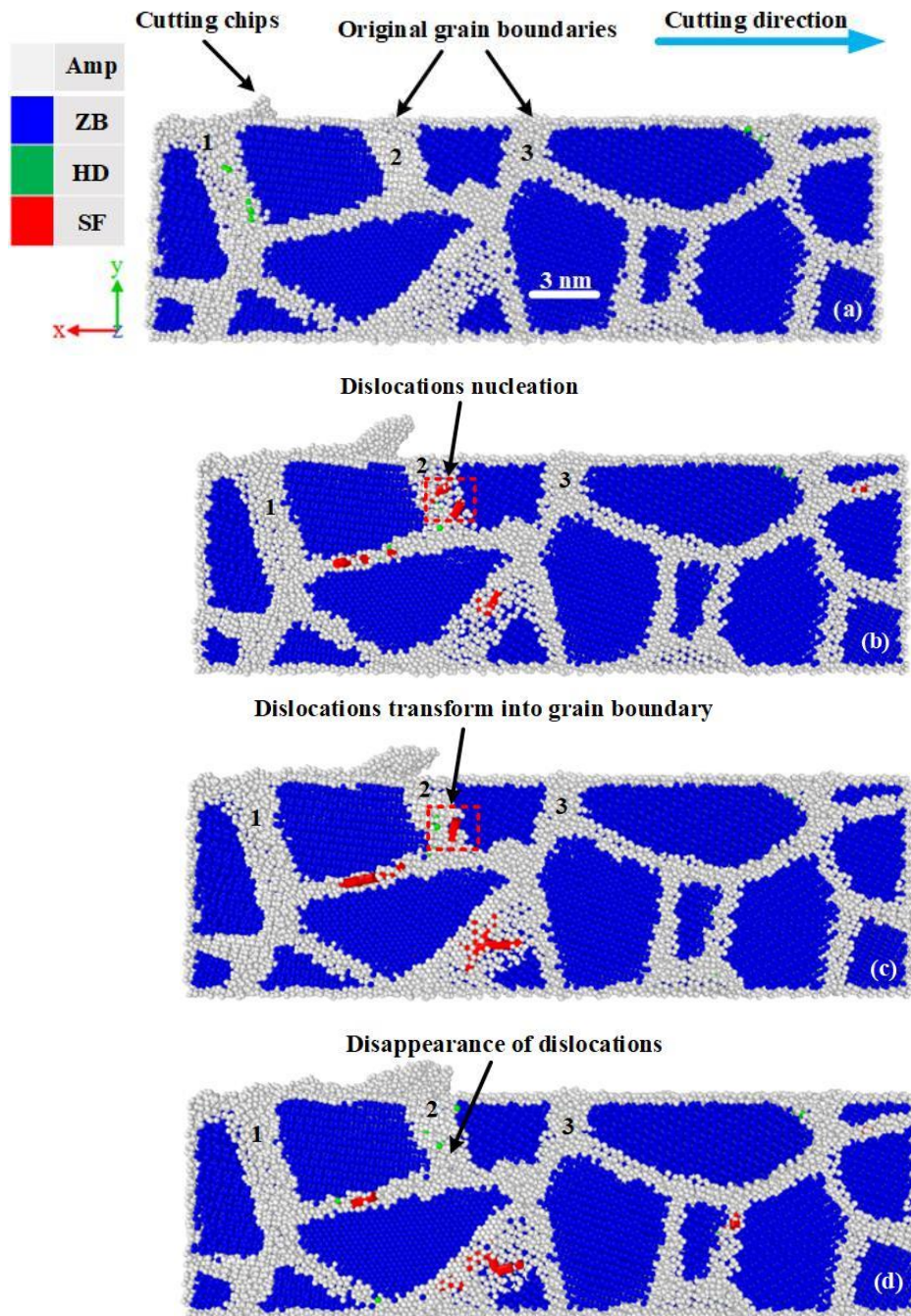
The choice of potential function can make a significant difference on the accuracy of MD results. It is important to choose a robust potential especially when it concerns studying aspects of fracture, wear and plasticity of a material. In this investigation, the cutting of GaAs with a diamond tool required describing the interactions between and among three types of atoms namely, Ga, As and C atoms. Due to the unavailability of a single many-body potential parameterized to describe all these atoms, a hybrid scheme was employed here in a hybrid/overlay scheme offered by LAMMPS. For the sake of brevity and avoiding repeated information, the details of the potential function (which is readily available from the respective papers) are not repeated here, but generally speaking, the covalently bonded interactions of C-C and the Ga-Ga, As-As and Ga-As interactions were all described by the analytical bond order potential developed by the research group of Albe et al. [27][28]. As for the cross interactions between the atoms of the diamond tool and the Gallium Arsenide workpiece (Ga-C and As-C), a Ziegler-Biersack-Littmark (ZBL) potential function [29] (`pair_style zbl` in LAMMPS) was used which simply requires the atomic number and cut off parameters as an input.

## **3 Results and discussions**

### *3.1 Microstructural changes*

Taking a test case of cutting depth of 2 nm and scratch velocity of 200 m/s, figure 3

shows a simulation output wherein blue color atoms, white color atoms, green color atoms and red color atoms represent the perfect zinc blende (ZB) structure, amorphous (Amp) structure, hexagonal diamond (HD) structure and stacking faults (SF) respectively.





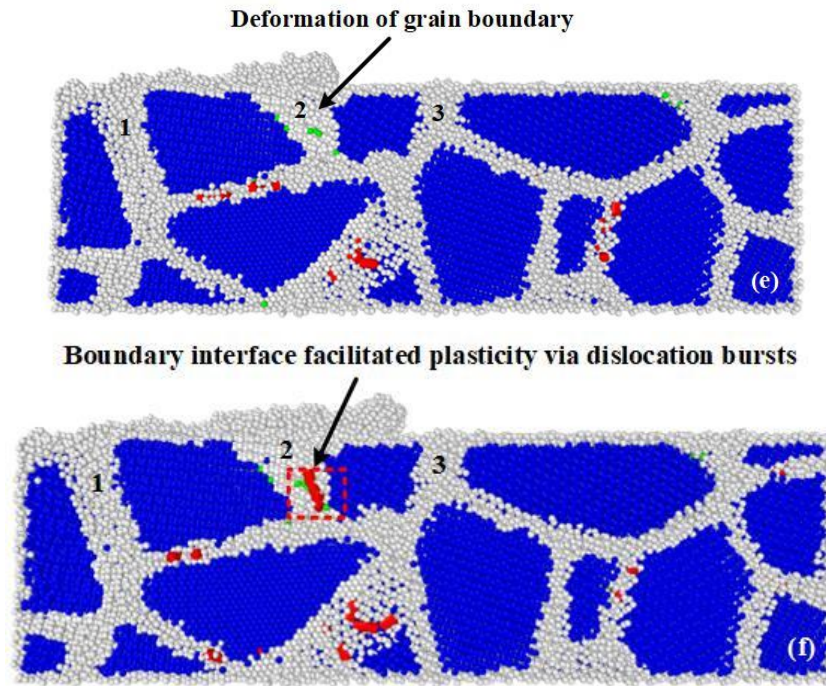
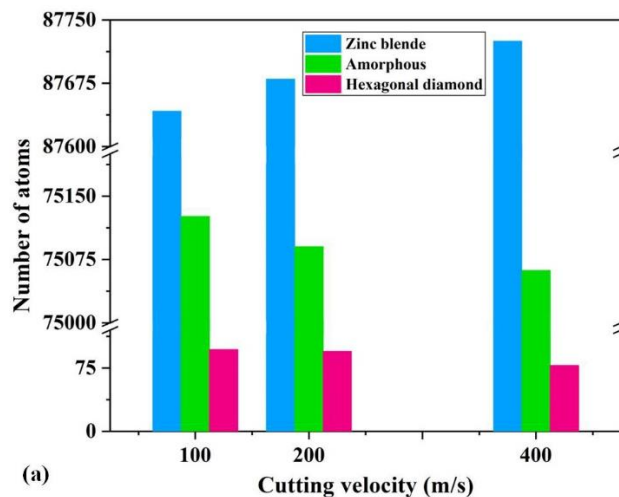


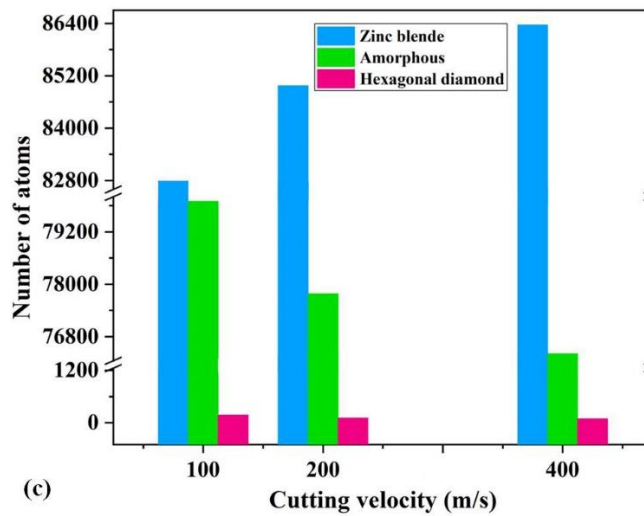
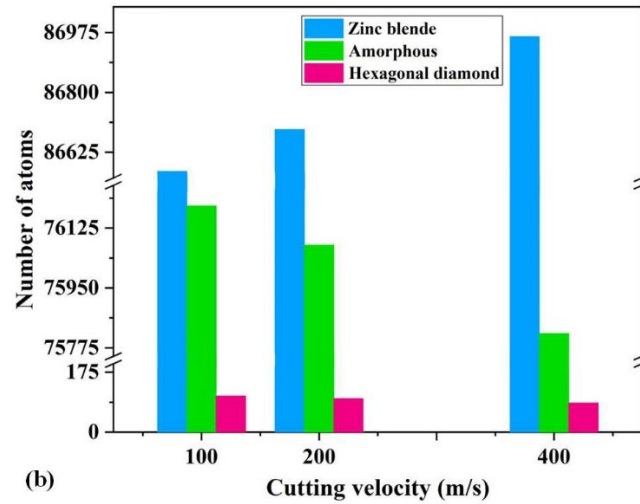
Fig. 3. Cross-sectional image of the polycrystalline GaAs (tool is kept hidden for visualization and cutting is performed at a depth of 2 nm and scratch velocity of 200 m/s). The snapshots are taken at cutting distances (a) 3 nm, (b) 6 nm, (c) 7 nm, (d) 8 nm, (e) 9 nm and (f) 12 nm. Pictures were processed using OVITO.

As shown in Fig. 3(a), the grain boundaries marked by 1, 2 and 3 were chosen as the sites of analysis for post-processing visualization of the dislocation and stacking fault structures using the second nearest neighbor scheme relying on an extended common neighbor analysis implemented in OVITO [30]. Fig. 3(b) highlights the initiation of the nucleation of dislocations at several places in the grain boundary 2 (GB 2). With subsequent tool travel, i.e. at the cutting distance of 7 nm, a part of the dislocations created at a cutting distance of 6 nm started to transform to grey color atoms and became a part of GB 2. Subsequently, the dislocations within GB 2 disappeared and transformed into a grain boundary marked by grey color atoms at the cutting distance of 8 nm shown

in Fig. 3(d). Meanwhile, a small number of green atoms were found present in the grain boundaries. At the cutting distance of 8 nm, the GB 2 widens and became thicker as may be seen in Fig. 3(e). Finally, as shown in Fig. 3(f), a large dislocation burst appeared in the GB 2 at the cutting distance of 12 nm, and a few grey color atoms transformed into red color atoms. The observation reported here remains consistent in all the simulation test cases.

Meanwhile, the evolution of structural changes in polycrystalline GaAs was quantified as a function of cutting velocity at various depths of scratch (see Fig. 4). It must be noted here that the ABOP potential function used in this study does not predict the energy differences between the cubic diamond and hexagonal closed packed diamond and hence the observation of hexagonal diamond in this study is a mere reflection on the faulted diamond cubic structure [31]. Overall, results in figure 4 show that a higher cutting velocity leads to a reduced degree of structural transformation in the material. Additionally, it was observed that all such microstructural transformations initiate preferentially in the grain boundaries. It implies that grain boundaries are preferred sites of nucleation of dislocations during the scratching of a polycrystalline substrate.



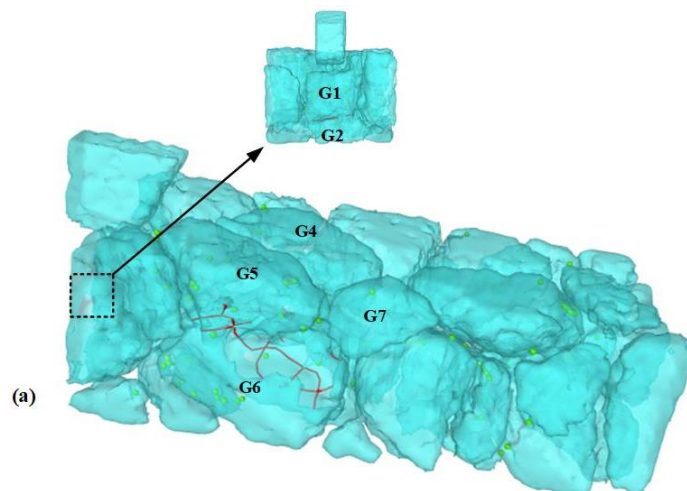


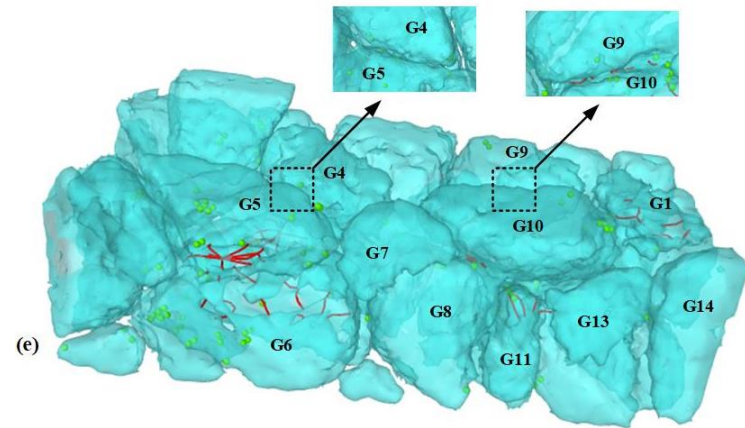
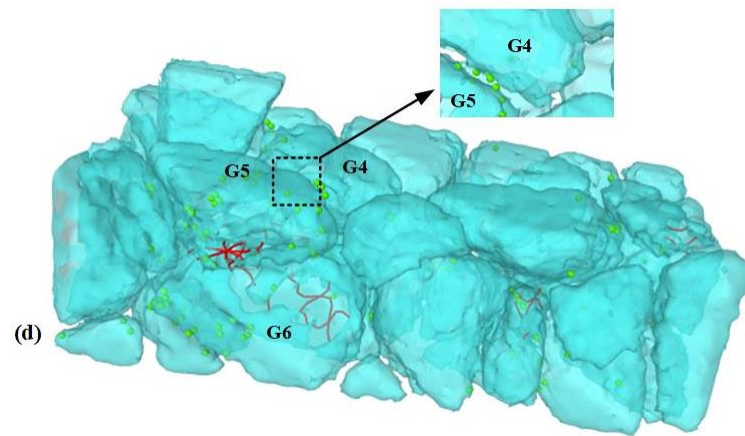
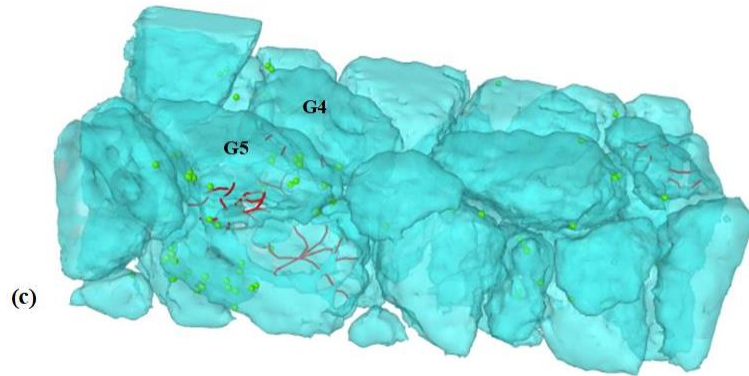
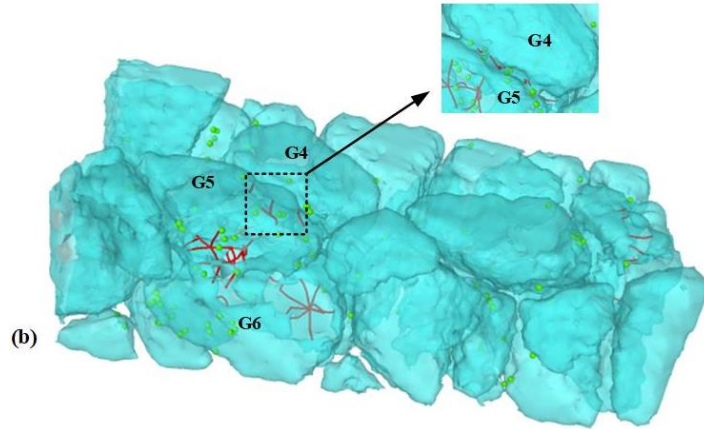
**Fig. 4.** The evolution of microstructure changes in polycrystalline GaAs during nanoscratching process with various cutting velocities under cutting depth of 0.5 nm (a), 1 nm (b) and 2 nm (c). Note here that the occurrence of the formation of hexagonal diamond is a mere artefact since the potential function used in this study does not distinguish energy differences between cubic and hexagonal phases.

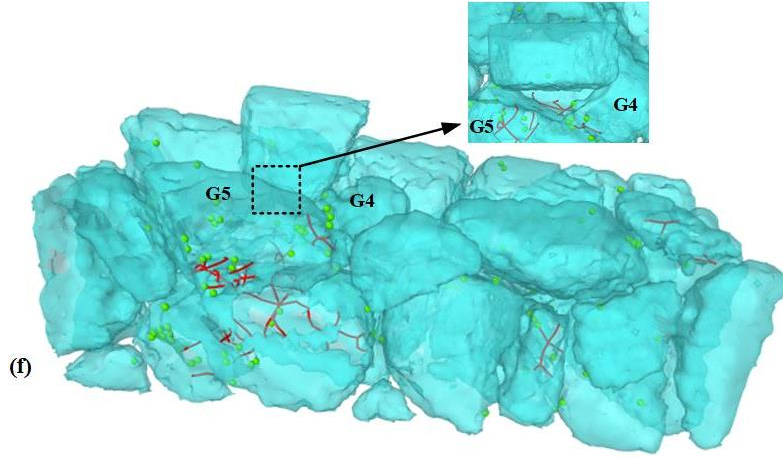
### 3.2 Analysis of dislocation nucleation

Fig. 5 shows the details of dislocation nucleation. According to the three-dimensional (3D) images, the nucleation of dislocations (marked by red color atoms) occurred in the grain boundaries, which is consistent with the two-dimensional (2D) images shown

earlier in Fig. 3. As no dislocations were found inside of the individual grains of polycrystalline GaAs, it indicated that the grain boundaries are softer than the grains and deform swiftly. This phenomenon is in accordance with the recently reported work on polycrystalline silicon carbide material [32]. Additionally, when the diamond tool penetrated the polycrystalline GaAs at 3 nm (see Fig. 5(a)), the two clusters of the dislocations were found in G1-G2 and G5-G6, respectively. Subsequently, the dislocation nucleation diffused through G4-G5 and the right corner of the polycrystalline GaAs workpiece, as illustrated in Fig. 5(b). In accordance with Fig. 5(c) and Fig. 5(d), the dislocation nucleation kept reappearing when the diamond tool passed through the edge of the grain boundary between G4 and G5. When the diamond tool started to penetrate the grain boundary between G4 and G5, there was no dislocation nucleation in G4-G5 (see figure 5(e)). The dislocation nucleation was distributed across the G5-G6 and G7-G8-G9-G10-G11-G13. Finally, the dislocation nucleation reoccurred in the G4-G5, as showed in Fig. 5(f), while the diamond tool cuts the grain boundary between G4 and G5.





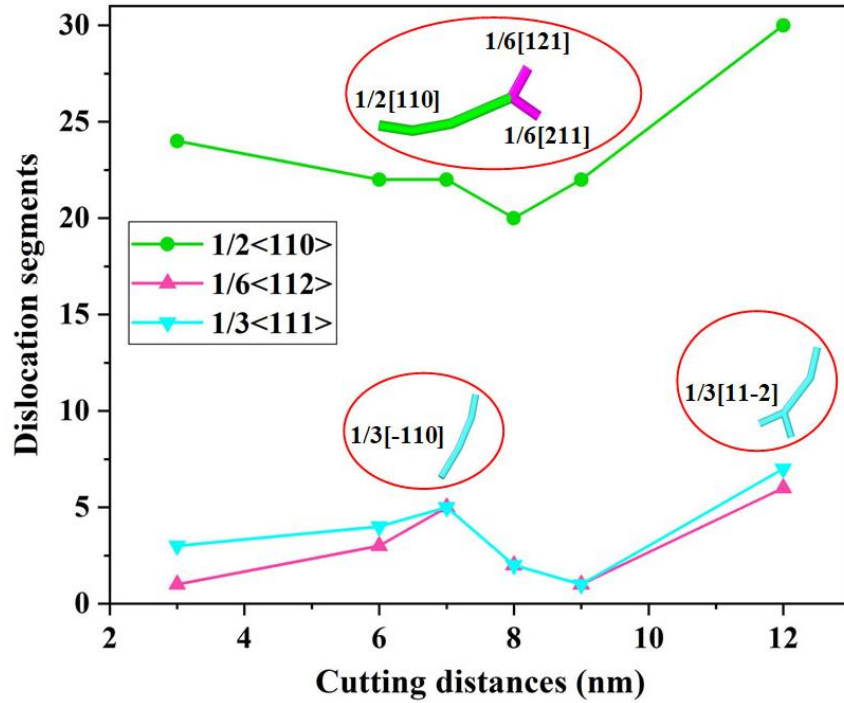


**Fig. 5.** The movement of dislocations in the polycrystalline GaAs at (a) 3 nm (b), 6 nm (c), 7 nm (d), 8 nm (e), 9 nm and (f) 12 nm.

In terms of the quantitative analysis, the number of dislocation segments extracted from the MD data is shown in Fig. 6. It can be seen that the presence of  $1/2\langle 110 \rangle$  type dislocations dominated others which was responsible for the incipient plasticity observed in the polycrystalline GaAs. The two other dislocations of type  $1/6\langle 112 \rangle$  and  $1/3\langle 111 \rangle$  were also present and were of the same length for the duration of cutting studied here. During the simulation, the dislocation with  $1/2[110]$  Burgers vector was observed to split into two Shockley partials with one having  $1/6[121]$  Burgers vector and the other having  $1/6[211]$  Burgers vector. The dissociation reaction can be represented as  $1/2[110] = 1/6[121] + 1/6[211]$ .

Also, the  $1/3\langle 111 \rangle$  dislocation with  $[-110]$  Burgers vector at 7 nm cutting distance appears to dissociate to a  $1/3\langle 111 \rangle$  dislocation with  $[11-2]$  Burgers vector. This phenomenon suggests the shuffle set dislocations could transit to glide set dislocations under large shear stress caused by the scratching tool [33]. The occurrence of the dual

slip mechanisms was seen an important factor driving plasticity in poly GaAs in sharp contrast to a single GaAs.



**Fig. 6.** Variation in the extent of dislocation segments and dislocations images with cutting distance.

An important physical quantity, dislocation density, was employed to describe the total length of dislocation lines contained in a unit volume of polycrystalline GaAs. The dislocation density was calculated by Eq. (1) [34].

$$\rho = \frac{L}{V} \quad (1)$$

where the  $L$  and  $V$  represent the total length of dislocation lines ( $\text{\AA}$ ) and volume of workpiece ( $\text{\AA}^3$ ), respectively. Consequently, the evolution of dislocation density as a functional of scratching distance is illustrated in Fig. 7.

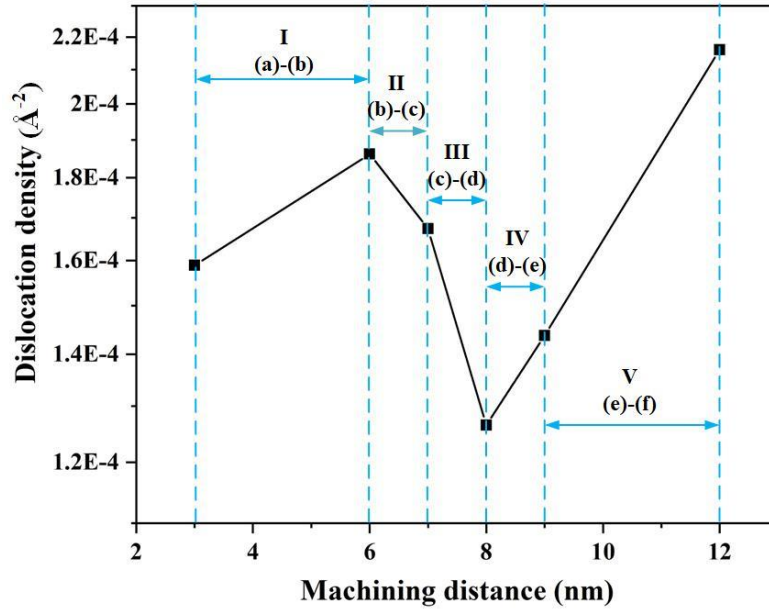


Fig. 7. Evolution of the dislocation density as a function of machining distance.

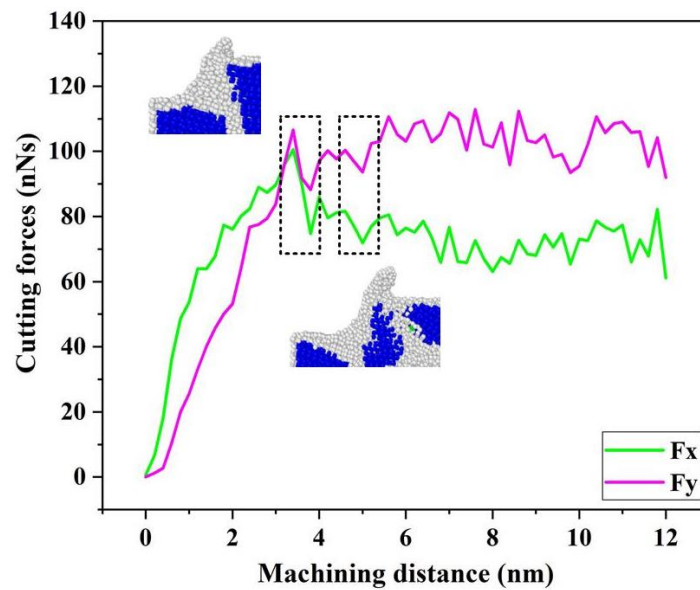
The dislocation density curve was seen to consist of five stage (I, II, III, IV and V). The dislocation density in the Ist stage was seen to increase which indicated the initiation of dislocation nucleation within the grain boundary. The dislocation density in the IInd and IIIrd stage decreases implied that certain dislocations transform to grain boundaries vis-a-vis disappearing of certain dislocations in a certain grain boundary as shown in Fig. 3 and Fig. 5. The dislocation density of the IV and V stage experienced a significant increase indicating that the diamond tool propagated through the grain boundaries to cause more dislocations and the cycle keeps repeating.

### 3.3 Differences in the cutting of single crystal and polycrystalline GaAs

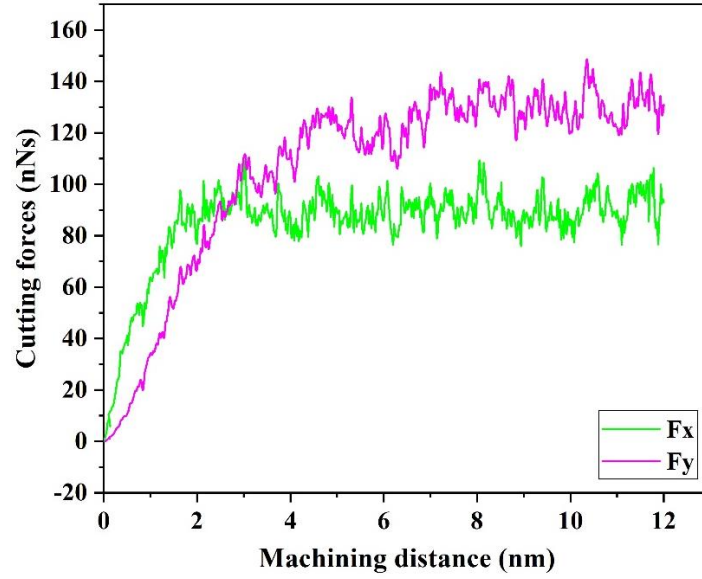
During this investigation, additional MD simulations were performed to benchmark the scratch forces namely, the lateral force ( $F_x$ ) and normal force ( $F_y$ ) during cutting of polycrystalline GaAs and single crystal GaAs. Fig 8 shows the evolution of the scratch



forces obtained from the MD simulations while cutting polycrystalline GaAs and single crystal GaAs substrates. Initially, until the onset of chip formation (unsteady cutting condition), the lateral force ( $F_x$ ) was seen to be larger than the normal force ( $F_y$ ) and once the machining achieved a steady-state, then the normal force ( $F_y$ ) becomes larger than the lateral force. In this study, under the same scratching condition (depth of cut of 2 nm and cutting velocity of 200 m/s), the lateral ( $F_x$ ) and normal force ( $F_y$ ) while cutting polycrystalline GaAs were about 70 nN and 110 nN respectively while the forces during cutting of the single crystal GaAs were of the order of 90 nN and 130 nN, respectively. The variation trend of MD normal force simulation can be validated qualitatively to some extent by our recent experimental results reported in [35].



(a) Scratch forces during cutting of a polycrystalline GaAs



(b) Scratch forces during cutting of a single crystal GaAs

**Fig. 8.** Evolution of cutting forces i.e. lateral ( $F_x$ ) and normal force ( $F_y$ ) at a cutting velocity of 200 m/s and at depth of cut of 2 nm.

Furthermore, by comparing Fig. 8(a) and Fig. 8(b), it can be seen that the lateral ( $F_x$ ) and normal force ( $F_y$ ) smoothly undulated from crests to troughs during cutting of polycrystalline GaAs. The reason for this is that the cutting force drops as the grains started to slide along an easy slip direction and when the grain boundary paved the way for the plastic deformation causing the cutting energy to be mainly concentrated in the grain boundaries. Beyond a certain threshold, the grain boundary collapses releasing a burst of deformation energy which leads to wave troughs of the cutting force. Additional calculations of the specific cutting energy ( $e_c$ ) and friction coefficient ( $F_x/F_y$ ) were also made. The specific cutting energy is defined as the work done by the tool in removing the unit volume of material and it can be calculated as [36].

$$e_c = \frac{R}{b \times t} \quad (2)$$

where  $R$  refers to the resultant force  $\sqrt{F_x^2 + F_y^2}$  while  $b$  and  $t$  represents the width of cut and depth of cut, respectively.

As shown in Table 2, the resultant cutting force, specific cutting energy and kinetic coefficient of friction values for cutting polycrystalline substrate were seen to be lower in magnitude compared to cutting single crystal GaAs.

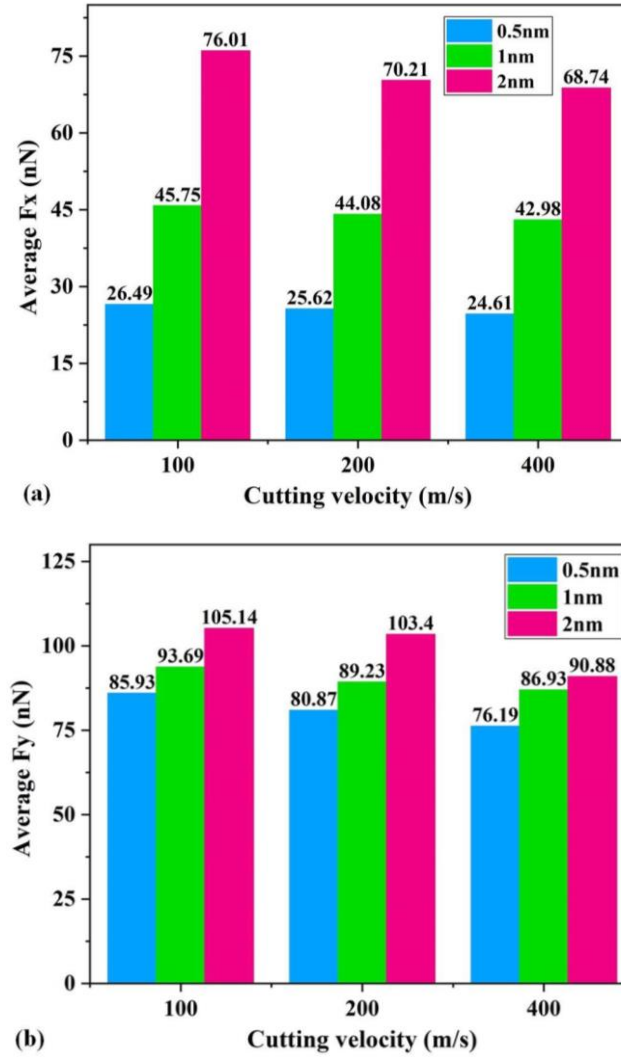
**Table 2:** Comparison of cutting results for single crystal GaAs and polycrystalline GaAs

Workpiece	$F_x$ (nN)	$F_y$ (nN)	$F_r$ (nN)	Specific cutting energy (GPa)	Friction coefficient
Single crystal GaAs	91.56	127.35	156.85	27.46	0.719
Polycrystalline GaAs	70.21	103.40	124.98	21.88	0.679

The machining force results indicated that the polycrystalline GaAs was more machinable than the single crystal GaAs. This is due to the presence of grain boundaries which eases the ductile deformation of a polycrystalline substrate.

### 3.4 Influence of cutting depth and cutting speed

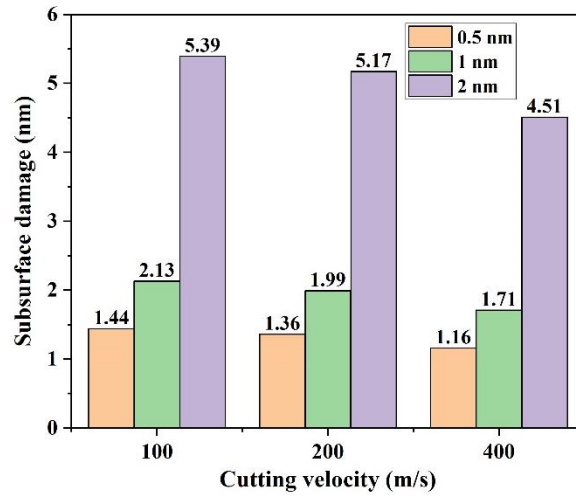
In this section, the influence of cutting depth and cutting speed on the cutting forces, sub-surface damage depth and the cutting temperature are reported. As shown in Fig. 9, both lateral ( $F_x$ ) and normal forces ( $F_y$ ) during cutting of the polycrystalline GaAs decreased with the increase of cutting speed or decreasing depth of cut. It was further observed that the normal force ( $F_y$ ) continues to be higher than the lateral force ( $F_x$ ) in all cases of scratching.



**Fig. 9.** The average value of the lateral ( $F_x$ ) and normal forces ( $F_y$ ) under different cutting velocities and depth of cut in cutting of polycrystalline GaAs.

Next, the sub-surface damage depth during cutting of polycrystalline GaAs was estimated as a function of different speeds and depth of cut which is shown in figure 10. It can be seen that the damage depth reduces with the increase of cutting speed which indicates that high strain rate applied during cutting decreases the sub-surface damage. A maximum sub-surface damage reduction of 16.32% could be achieved while cutting at 400 m/s at a depth of cut of 2 nm in comparison to cutting at 100 m/s at the

same depth of cut.



**Fig. 10.** Sub-surface damage depth at different cutting velocities and depth of cut

Finally, the temperature variation as a function of depth of cut and cutting speed was estimated and shown in figure 11. Higher speed of cutting and higher depths of cutting were both seen to accompany an increase in the cutting temperature in the plastic zone. The combined information of the temperature and stresses acting in the cutting zone could be used as a vital information to predict the microstructural changes in the cutting zone and we shall expand on this aspect in our future work.

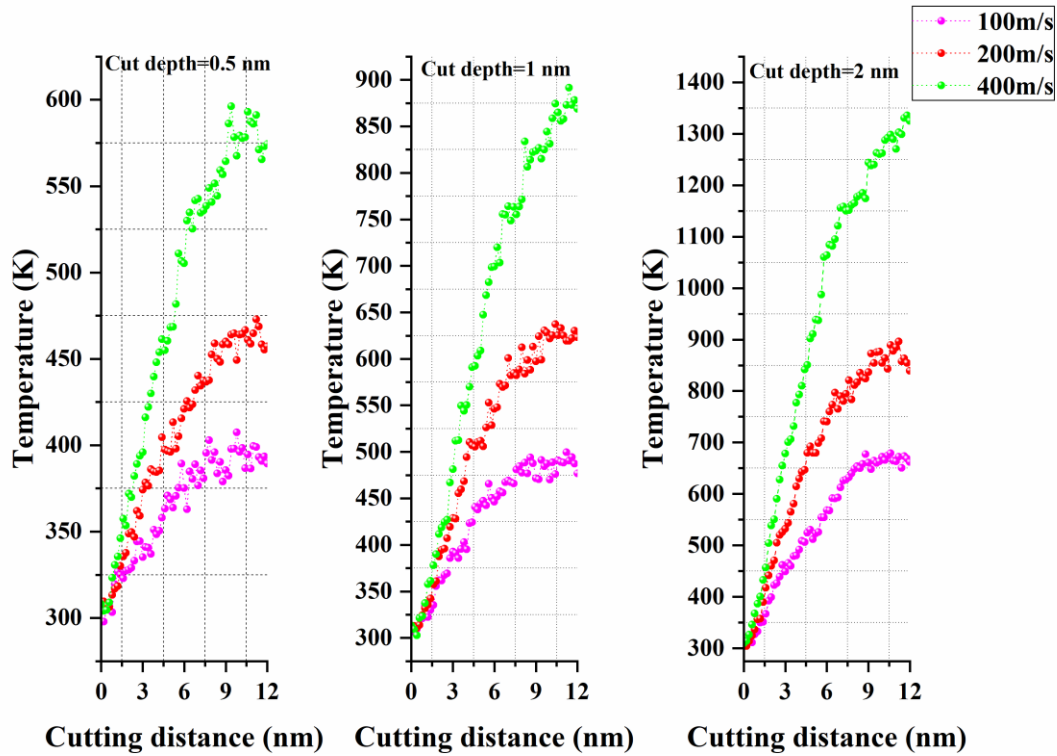


Fig. 11. Peak temperature variation at various scratch speeds and depth of cuts

## 4 Conclusions

In this work, the deformation mechanism of polycrystalline GaAs during nanoscratching was investigated by the MD simulations and benchmarked to single crystal GaAs. During the simulations, the scratch depth, speed of scratching (thus the applied strain rate) and microstructure of the workpiece (polycrystalline vs single crystal GaAs) were varied and output parameters such as the scratch forces (and specific cutting energy), kinetic coefficient of friction, cutting temperature, sub-surface damage and dislocation structures were extracted and analysed. In light of these extracted parameters and the analysis performed, the following conclusions can be drawn:

1. The presence of grain boundaries eases the deformation of the polycrystalline GaAs as opposed to single crystal GaAs. It was discovered that the grain

boundaries can become the incipient sites of dislocation nucleation and thus become the weak links in a polycrystal as opposed to a single crystal which has no such weaker links. The ease of plastic deformation of the grain boundaries compared to the individual grains makes polycrystals more easily deformed than the single crystals.

2. The cutting forces showed a unique cyclic wave crest to wave troughs transition while cutting polycrystalline GaAs in contrast to the cutting of the single crystal GaAs. This was attributable to the periodic arrest of the dislocations in the grain boundaries followed by collapsed grain boundaries as a result of the continuous tool scratching.
3. The friction coefficient and the specific cutting energy were found to be higher for scratching single crystal GaAs than for polycrystalline GaAs and also the normal scratch force achieves a higher magnitude over the lateral scratch force once the scratching has achieved a steady state.
4. Scratch forces and the sub-surface damage were observed to reduce with an increase in the scratch velocity and to increase with the increasing depth of scratch. However, the cutting temperature increases with the increase in scratch speed and the scratch depth.
5. The  $\frac{1}{2}\langle 110 \rangle$  was found to be the main type of dislocation responsible for ductile plasticity in polycrystalline GaAs which splits into Shockley partials connected by an Internal Stacking Fault (ISF) leading to dissociation of the parent dislocation in  $\frac{1}{6}\langle 121 \rangle$  and  $\frac{1}{6}\langle 211 \rangle$  type dislocations.

## **Acknowledgement**

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## **Data statement**

All data underpinning this publication are openly available from the University of Strathclyde Knowledge Base at <https://doi.org/10.15129/ed05ea22-ab4c-4dd7-af36-4dfb91bd1cef>.

**Conflict of Interest:** The authors declare that they have no conflict of interest.



## References

- [1] H. J. Ahn, W. Il Chang, S. M. Kim, B. J. Park, J. M. Yook, and Y. S. Eo, “28 GHz GaAs pHEMT MMICs and RF front-end module for 5G communication systems,” *Microw. Opt. Technol. Lett.*, vol. 61, no. 4, pp. 878–882, 2019.
- [2] T. P. Chow, I. Omura, M. Higashiwaki, H. Kawarada, and V. Pala, “Smart power devices and ICs using GaAs and wide and extreme bandgap semiconductors,” *IEEE Trans. Electron Devices*, vol. 64, no. 3, pp. 856–873, 2017.
- [3] A. N. Gulluoglu and C. T. Tsai, “Dislocation generation in GaAs crystals grown by the vertical gradient freeze method,” *J. Mater. Process. Technol.*, vol. 102, no. 1, pp. 179–187, 2000.
- [4] F. M. Kiessling, M. Albrecht, K. Irmscher, R. Krause-Rehberg, W. Ulrici, and P. Rudolph, “Defect distribution in boron-reduced GaAs crystals grown by vapour-pressure-controlled Czochralski technique,” *J. Cryst. Growth*, vol. 310, no. 7–9, pp. 1418–1423, 2008.
- [5] R. K. Ahrenkiel, S. W. Johnston, B. M. Keyes, and D. J. Friedman, “Transport properties of GaAs<sub>1-x</sub>N<sub>x</sub> thin films grown by metalorganic chemical vapor deposition,” *Appl. Phys. Lett.*, vol. 77, no. 23, pp. 3794–3796, 2000.
- [6] S. Zhou, L. Ai, M. Qi, S. Wang, A. Xu, and Q. Guo, “Bi-induced highly n-type carbon-doped InGaAsBi films grown by molecular beam epitaxy,” *J. Mater. Sci.*, vol. 53, no. 5, pp. 3537–3543, 2018.
- [7] S. Blin *et al.*, “Wireless communication at 310 GHz using GaAs high-electron-

- mobility transistors for detection,” *J. Commun. Networks*, vol. 15, no. 6, pp. 559–568, 2013.
- [8] K. Alberi, B. Fluegel, H. Moutinho, R. G. Dhere, J. V Li, and A. Mascarenhas, “Measuring long-range carrier diffusion across multiple grains in polycrystalline semiconductors by photoluminescence imaging,” *Nat. Commun.*, pp. 1–7, 2013.
- [9] D. M. Wilt, M. A. Smith, W. Maurer, D. Scheiman, and P. P. Jenkins, “GaAs photovoltaics on polycrystalline Ge substrates,” *Conf. Rec. 2006 IEEE 4th World Conf. Photovolt. Energy Conversion, WCPEC-4*, vol. 2, no. June 2006, pp. 1891–1894, 2006.
- [10] M. Yamaguchi and Y. Itoh, “Efficiency considerations for polycrystalline GaAs thin-film solar cells,” *J. Appl. Phys.*, vol. 60, no. 1, pp. 413–417, 1986.
- [11] M. K. Sharma and D. P. Joshi, “Electrical conduction model for polycrystalline GaAs films,” *J. Appl. Phys.*, vol. 102, no. 3, pp. 1–8, 2007.
- [12] M. Imaizumi *et al.*, “Low-temperature growth of GaAs polycrystalline films on glass substrates for space solar cell application,” *J. Cryst. Growth*, vol. 221, no. 1–4, pp. 688–692, 2000.
- [13] J. C. Bourgoin, “Polycrystalline GaAs for large area imaging detectors,” *Nucl. Instruments Methods Phys. Res. Sect. A Accel. Spectrometers, Detect. Assoc. Equip.*, vol. 466, no. 1, pp. 9–13, 2001.
- [14] J. H. Epple, K. L. Chang, C. F. Xu, G. W. Pickrell, K. Y. Cheng, and K. C. Hsieh, “Formation of highly conductive polycrystalline GaAs from annealed

- amorphous (Ga,As),” *J. Appl. Phys.*, vol. 93, no. 9, pp. 5331–5336, 2003.
- [15] J. D. Song, W. J. Choi, J. I. Lee, J. M. Kim, K. S. Chang, and Y. T. Lee, “Optical and structural properties of InGaAs/InP double quantum wells grown by molecular beam epitaxy with polycrystalline GaAs and GaP decomposition sources,” *Phys. E Low-Dimensional Syst. Nanostructures*, vol. 32, no. 1-2 SPEC. ISS., pp. 234–236, 2006.
- [16] S. Z. Chavoshi, S. Xu, and S. Goel, “Addressing the discrepancy of finding the equilibrium melting point of silicon using molecular dynamics simulations,” *Proc. R. Soc. A Math. Phys. Eng. Sci.*, vol. 473, no. 2202, pp. 1–9, 2017.
- [17] S. Goel, N. Haque Faisal, X. Luo, J. Yan, and A. Agrawal, “Nanoindentation of polysilicon and single crystal silicon: Molecular dynamics simulation and experimental validation,” *J. Phys. D. Appl. Phys.*, vol. 47, no. 27, 2014.
- [18] S. Goel, J. Yan, X. Luo, and A. Agrawal, “Incipient plasticity in 4H-SiC during quasistatic nanoindentation,” *J. Mech. Behav. Biomed. Mater.*, vol. 34, pp. 330–337, 2014.
- [19] S. Goel *et al.*, “Horizons of modern molecular dynamics simulation in digitalized solid freeform fabrication with advanced materials,” *Mater. Today Chem.*, vol. 18, p. 100356, 2020.
- [20] P. Hirel, “Atomsk: A tool for manipulating and converting atomic data files,” *Comput. Phys. Commun.*, vol. 197, pp. 212–219, 2015.
- [21] G. Voronoi, “Nouvelles applications des paramètres continus à la théorie des formes quadratiques. Deuxième mémoire. Recherches sur les paralléloèdres

- primitifs,” *J. für die reine und Angew. Math.*, vol. 1908, no. 134, pp. 198–208, 2009.
- [22] G. Lejeune Dirichlet, “Über die Reduction der positiven quadratischen Formen mit drei unbestimmten ganzen Zahlen.,” *J. für die reine und Angew. Math.*, vol. 40, pp. 209–227, 1850.
- [23] S. Nose, “A unified formulation of the constant temperature molecular-dynamics methods,” *J. Chem. Phys.*, vol. 81, pp. 511–519, 1984.
- [24] S. J. Plimpton, “Fast parallel algorithms for short range molecular dynamics,” *J. Comput. Phys.*, vol. 117, pp. 1–19, 1995.
- [25] D. T. Infield *et al.*, “Main-chain mutagenesis reveals intrahelical coupling in an ion channel voltage-sensor,” *Nat. Commun.*, vol. 9, no. 1, pp. 1–10, 2018.
- [26] A. Stukowski, “Visualization and analysis of atomistic simulation data with OVITO-the Open Visualization Tool,” *Model. Simul. Mater. Sci. Eng.*, vol. 18, no. 1, 2010.
- [27] P. Erhart and K. Albe, “Analytical potential for atomistic simulations of silicon, carbon, and silicon carbide,” *Phys. Rev. B - Condens. Matter Mater. Phys.*, vol. 71, no. 3, pp. 1–14, 2005.
- [28] K. Albe, K. Nordlund, J. Nord, and A. Kuronen, “Modeling of compound semiconductors: Analytical bond-order potential for Ga, As, and GaAs,” *Phys. Rev. B - Condens. Matter Mater. Phys.*, vol. 66, no. 3, pp. 352051–3520514, 2002.
- [29] J. F. Ziegler, M. D. Ziegler, and J. P. Biersack, “SRIM - The stopping and

- range of ions in matter (2010),” *Nucl. Instruments Methods Phys. Res. Sect. B Beam Interact. with Mater. Atoms*, vol. 268, no. 11–12, pp. 1818–1823, 2010.
- [30] E. Maras, O. Trushin, A. Stukowski, T. Ala-Nissila, and H. Jónsson, “Global transition path search for dislocation formation in Ge on Si(001),” *Comput. Phys. Commun.*, vol. 205, pp. 13–21, 2016.
- [31] S. Goel and A. Stukowski, “Comment on ‘incipient plasticity of diamond during nanoindentation’ by C. Xu, C. Liu and H. Wang, : RSC Advances, 2017, 7, 36093,” *RSC Adv.*, vol. 8, no. 10, pp. 5136–5137, 2018.
- [32] S. Goel, X. Luo, P. Comley, R. L. Reuben, and A. Cox, “Brittle-ductile transition during diamond turning of single crystal silicon carbide,” *Int. J. Mach. Tools Manuf.*, vol. 65, pp. 15–21, 2013.
- [33] Z. Li and R. C. Picu, “Shuffle-glide dislocation transformation in Si,” *J. Appl. Phys.*, vol. 113, no. 8, pp. 1–7, 2013.
- [34] S. H. He, B. B. He, K. Y. Zhu, and M. X. Huang, “Evolution of dislocation density in bainitic steel: Modeling and experiments,” *Acta Mater.*, vol. 149, pp. 46–56, 2018.
- [35] J. Chen, X. Luo, F. Ding, X. Rao, and J. Zhang, “Fundamental study of diamond turning of single crystal gallium arsenide,” *Precis. Eng.*, vol. 62, pp. 71–82, 2020.
- [36] S. Goel, “The current understanding on the diamond machining of silicon carbide,” *J. Phys. D. Appl. Phys.*, vol. 47, no. 24, 2014.