# 1 Title: Elytra coupling of the ladybird *Coccinella septempunctata* functions as an

# 2 energy absorber in intentional falls

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#### 20 Abstract:

21 Some insects, such as bees, wasps, and bugs, have specialized coupling structures 22 to synchronize the wing motions in flight. Some others, such as ladybirds, are equipped 23 with coupling structures that work only at rest. By locking elytra into each other, such 24 structures provide hindwings with a protective cover to prevent contamination. Here, 25 we show that the coupling may play another significant role: contributing to energy 26 absorption in falls, thereby protecting the abdomen against mechanical damage. In this 27 combined experimental, numerical and theoretical study, we investigated free falls of 28 ladybirds (Coccinella septempunctata), and discovered that upon collision to the 29 ground, the coupling may fail and the elytra may unlock. This unlocking of the coupling 30 increased the energy absorption by 33%, in comparison to when the elytra remain 31 coupled. Using micro-CT scanning, we developed comparative models that enabled us 32 to simulate impact scenarios numerically. Our results showed that unlocking of the 33 coupling, here called elytra splitting, reduces both the peak impact force and rebound 34 velocity. We fabricated the insect-inspired coupling mechanism using 3D printing and 35 demonstrated its application as a damage prevention system for quadcopters in 36 accidental collisions.

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38 Keyword: ladybird, elytra coupling, landing, energy absorption.

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#### 40 **1. Introduction**

41 The performance of energy absorbing structures in vertebrates, such as turtle shells 42 Terrapene carolina (Rhee et al., 2009; Wu et al., 2017), sheep horns Ovis canadensis 43 (Fuller and Donahue, 2021), and woodpecker beaks *Melanerpes carolinus* (Lee *et al.*, 44 2014), has fascinated researchers for many decades (Tarakanova and Buehler, 2012). 45 For instance, woodpeckers perform repeated strikes at a speed of 6-7 m/s while pecking 46 for food, resulting in an impact deceleration of the order of 1000 g (May et al., 1979). 47 Horns of bighorn sheep, as another example, resist forces of up to 3400 N during fights 48 (Trim *et al.*, 2011). The massive impact loads experienced by such biological structures 49 and the absence of any catastrophic failure, suggest that they have undergo adaptations

to excessive impact stresses (Drake et al., 2016). Similar to vertebrates, energy 50 51 absorbing structures play an important role in the survival of invertebrates. Orb webs 52 of the spider Argiope aurantia, which captures insects in midflight, can absorb prey 53 kinetic energy without breaking (Foelix, 2011; Sensenig et al., 2012). Shells of the 54 mantis Odontodactylus scyllarus withstand thousands of high-velocity blows that the 55 animal delivers to its prey (Grunenfelder *et al.*, 2014). The leg muscles of the cockroach 56 Blaberus discoidalis not only generate the power for locomotion, but operate as active dampers that absorb energy during running (Full et al., 1998). 57

58 Beetles, one of the most diverse orders of insects with about 400 thousand species, represent striking examples of energy absorbing structures (Phan and Park, 2020). Their 59 60 elytra, which are in fact hardened forewings, are ventrally jointed to the body. By keeping the elytra closed, coupling structures enable the beetles to pass through 61 62 restricted environments without damaging the abdomen or hindwings (Dai and Yang, 63 2010; Breed and Ball, 1908; Dai et al., 2008; Frantsevich et al., 2005). Furthermore, when disturbed by predators, such as birds, amphibians, and reptiles, beetles play dead 64 and fall to the ground. During this displaying behavior, the elytra show a high energy 65 66 absorption capacity (Zhang et al., 2021; Burgio, 2013).

67 Knowing the internal morphology and material properties of elytra can provide 68 insights for the development of novel energy absorbing structures (Hao and Du, 2018), 69 a reason that has sparked research attention to the biomechanical properties of beetle elytra (Sun and Bhushan, 2012). Previous studies have shown that beetle elytra are 70 71 layered structures, primarily made of chitin microfibrils and protein (Jalali and 72 Heshmati, 2016; Tasdemirci et al., 2015), and consist of a dorsal and a ventral layer that 73 are interconnected by columnar trabecula (Kundanati et al., 2018; Hadley, 1986). With 74 the support of these structures and existing cavities, the elytra provide high strength and 75 efficient energy absorption that prevent damage (Hong et al., 2003; Gunderson and 76 Schiavone, 1995). Previous research, however, has fully neglected the potential role of 77 the coupling structure in the energy absorption properties of the elytra. To understand whether this might be the case, here we investigated the falling dynamics of the ladybird 78 79 Coccinella septempunctata and found that, while landing on the elytra, the coupling

structures can provide energy absorption and reduce both the rebound height and velocity after collision. We further show that the elytra coupling can offer novel insectinspired designs for mitigating body injuries as well as other applications that involve collisions.

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#### 85 2. Materials and Methods

## 86 2.1. Speciments

87 Ladybirds (C. septempunctata) were collected from Zhengzhou, Henan Province, 88 China (34.16°N, 112.42°E), and kept in a terrarium filled with indoor horse bean 89 seedlings (Vicia faba L.) in our laboratory in Guangzhou, Guangdong Province, China 90 (23.13°N, 113.27°E) at 26°C and 60% humidity with controlled day-night cycle 91 (12h/12h) (figure 1(a)). All ladybirds were fed with aphids (Megoura japonica) living 92 on the horse bean seedlings (Frazer and Gilbert, 1976; Zhang et al., 2021). For the 93 experiments, we randomly selected twenty ladybirds and recorded their body masses m. 94 The length a, width b, and height c of each sample were measured using vernier calipers 95 (AIRAJ 1331, AIRAJ, USA) (figure 1(b)).

#### 96 2.2. Calculation of energy absorption in falls

97 To measure the energy absorption during impacts, we dropped the insects (no 98 initial velocity) onto a hard surface through a transparent cylindrical tube (50 cm in 99 length and 4.5 cm in diameter) to eliminate the potential effect of the surrounding air 100 (figure 1(c)). To describe the orientation of the ladybirds, we set up a coordinate system 101 at the center of the abdominal plane o. The x axis pointed from point o towards the head, 102 and the z axis coincided with the normal direction of the abdomen. The y axis was 103 determined by the right-hand rule. The motion of the falling ladybirds was quantified 104 by roll, pitch, and yaw angles, about the axes x, y, and z, respectively. Two high-speed 105 cameras (Phantom, VEO-E 310L, USA), equipped with micro-lenses (Canon, 106 EF100mmf/2.8LISUSM, Japan), were used to record the falling ladybirds at 1000 fps. 107 The first high-speed camera (Camera #1) was used to observe the deformation of the 108 elytra at the impact. Using the second high-speed camera (Camera #2), we analyzed the

dynamics of falling, landing, and rebound before and after the impact, including the heights ( $h_1$ : before impact;  $h_2$ : after impact), translational velocities ( $v_1$ : before impact;  $v_2$ : after impact), and angular velocities ( $\omega_1$ : before impact;  $\omega_2$ : before impact) of the specimens.

113 Mechanical energy  $E_i$  of an object of mass m, height h, translational velocity 114 v, and angular velocity  $\omega$  can be obtained using the below equation:

$$E_i = mgh_i + \frac{1}{2} \left( mv_i^2 + J\omega_i^2 \right) \tag{1}$$

where, *J* is rotational inertia. The index *i* is the sequence of the record; i=1represents the initiation of the fall, and i=2 is when the rebound after collision with the ground has taken place.



**Figure 1.** Experimental setups for fall tests. (a) Ladybirds were kept in a terrarium filled with horse bean seedlings. (b) Measured dimensions of a ladybird body. (c) High-speed filming for the analysis of the falls and collisions of ladybirds. The defined *o-xyz* coordinates is shown here.

118 2.3. Impact force measurement

We built an experimental set-up with a force plate (HE6X6, AMTI, USA) to measure the impact force of ladybirds in free falls (figure 2(a)). The resolution of this force plate is 8.8 mN. To eliminate the influence of airflows and external vibrations on the measurements, we designed a screen shield made up of acrylic sheet around the force plate, and placed a foam board under the force plate. The sampling frequency and duration of the force plate were set as 100 Hz and 10 s, respectively. Additionally, a decibel meter was placed next to the force plate to measure the noise from the impacts. For experiments, we dropped the ladybirds (n = 20) from a constant height of 50 cm on the force plate with no initial velocity. This height corresponds to the average height of the plants on which ladybirds were observed.

# 129 2.4. Scanning electron microscopy

130 Scanning electron microscopy (SEM) was used to examine the microstructure of 131 the coupling between the elytra of the ladybirds. The fresh samples of elytra (n = 3)132 were fixed by soaking in 2.50% glutaraldehyde for 3 h at 26°C. The samples were then 133 cleaned with 0.10 mol/L phosphate buffer (pH = 7) for 20 minutes, and dehydrated 134 though an ascending ethanol series (at 75%, 80%, 85%, 90%, 95%, and 100% for 12 h) 135 before freeze drying. The samples were mounted on SEM stubs with graphite adhesive 136 tape. Coated in gold palladium, and observed under a SEM (FEI Quanta 200, Czech 137 Republic) in high-vacuum mode at 15 kV.

## 138 **2.5.** *Atomic force microscopy*

139 Atomic force microscopy (AFM) was used to characterize the stiffness of the 140 ladybird elytra. For this purpose, we used a Dimension Icon AFM (FastScan, BRUKER, 141 USA) (figure 2(b)) equipped with a probe (RTESPA-150, BRUKER, USA). Small 142 pieces (~ 1 mm  $\times$  1mm) were cut from the fresh elytra as specimens and covered by 143 wet cotton to prevent dehydration. We tested the specimens (n = 3) at ten sampling 144 points along the length of each specimen. In the peak force tapping mode, the vibrating 145 tip performed vertical indentation to record force curves for each sampling point on the 146 surface of the specimens. By real-time analysis of the force curves, the Young's 147 modulus of the specimens at each sampling point was obtained using the Derjaguin, 148 Muller, Toporov (DMT) model (Li et al., 2019).

#### 149 **2.6.** *Micro-computed tomography*

The dynamic X-ray micro-computed tomography (micro-CT) system from Shanghai Synchrotron Radiation Facility (SSRF) was employed to reconstruct the three-dimensional (3D) shape of the elytra of ladybirds (figure 2(c)). The beamline BL13W1 with 8.0 keV to 72.5 keV and several sets of X-ray imaging detectors with varying pixel sizes (0.19  $\mu$ m-24  $\mu$ m) were used to realize X-ray micro-CT in-line imaging (Wang *et al.*, 2020). In this experiment, the photon energy and pixel size were set as 15 keV and 3.25  $\mu$ m, respectively (Shang *et al.*, 2020; Xie *et al.*, 2020). A ladybird specimen, placed in a plastic tube, was mounted on the beamline positioner. After imaging, we processed the obtained images using the PITRE and Mimics software (Materialise, Belgium) for phase-sensitive X-ray image processing and tomography reconstruction to generate the 3D models of the elytra (Chen *et al.*, 2012).



**Figure 2.** Schematic diagram of experimental testing and imaging platform. (a) Diagram of impact force measurement platform. (b) Schematic diagram of the atomic force microscopy. (c) Schematic of the micro-computed tomography.

#### 161 **2.7. Finite element modeling**

162 Considering that the anterior parts of the elytra were always locked under the 163 impact, we focused on the posterior parts (x < 0) and defined two comparative models, 164 namely continuum and split models. We used our models to uncover how the coupling 165 can influence the energy absorption by the elytra in impacts. The first model was a 166 continuum shell that was geometrically similar to the profiles obtained from micro-CT 167 scans. This model was used to simulate the state at which the elytra are tightly coupled 168 with each other under the impact. By contrast, the second model consisted of two halves 169 of shells coupled with each other by a coupling that can be split under the impact.

For the dynamics of the models in falls, the governing equations were deployed, including equilibrium equation  $\sigma_{ij,j} + f_i = 0$ , geometric equation  $\varepsilon_{i,j} = (u_{i,j} + u_{j,i})/2$ ,

172 and constitutive equation  $\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2G\varepsilon_{ij}$ , in which  $\lambda = \frac{Ev}{(1+v)(1-2v)}$ , and

173  $G = \frac{E}{1+v}$ . Here, E and v represented elasticity modulus and Poisson's ratio, 174 respectively.

Based on the Hamilton variational principle, the governing equations for dynamiccollisions can be obtained (Hien and Kleiber, 1990).

$$M_i \partial \mathcal{K} + C_i \partial \mathcal{K} + K_i U = 0 \tag{2}$$

where,  $M_i$  was the mass matrix,  $C_i$  was the damping matrix,  $K_i$  was the stiffness matrix. Here,  $U^{(k)}$ ,  $U^{(k)}$ , and U were the acceleration, velocity and displacement vectors, respectively.

From the statistical analysis of the recorded videos of the ladybirds in falls, we found that different parts of the elytra might come in contact with the ground. This may result in different energy absorption capacities. To include the effect of this in our models, we defined the landing region by the landing angle, which is acute angle between the plane *xoy* and the ground. Here, we set three representative landing angles of  $0^{\circ}$ ,  $30^{\circ}$ , and  $60^{\circ}$ , respectively. We also used the material properties of the models as shown in Table 1 (Rivera *et al.*, 2020).

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#### Table 1. The material properties of the models

Elasticity modulus (GPa)	Poisson's ratio	Density (kg/m <sup>3</sup> )	Initial velocity (m/s)
3	0.3	500	2.75

188 During collisions, both models can absorb energy through the elastic deformation, 189 which convert the energy into the internal energy  $E_{in}$ . The internal energy of the models 190 could be written as:

$$E_{\rm in} = \sum_{i=1}^{n} \int_{v_i} f\left(\varepsilon_i\right) dv_i \tag{3}$$

191 where,  $f(\varepsilon_i)$  and  $v_i$  are the strain energy density and the volume of an element of 192 the model, and n is the total number of elements. 193

#### 194 **3. Results and discussion**

#### 195 **3.1. Energy absorption in falls**

196 The ladybirds (*C. septempunctata*,  $m = 27.04 \pm 6.62$  mg, n = 20) that were falling 197 from a height of 50 cm could adjust their orientations during falling, enabling them to 198 land on the elytra in 80% of the cases, regardless of the initial orientation. In all events, 199 we could measure the sound of 60-70 dB caused by the collision of the elytra to the 200 ground. This sound, which is analogous to the noise generated by a car running at 65 201 km/h, may reflect the intensity of the collisions (Behzad et al., 2007). However, the 202 ladybirds were not crashed to death by the impact force, which may suggest the 203 presence of an energy absorbing mechanism within the elytra. Considering that the part 204 at which the elytra collide to the ground may affect the energy absorption, we divided 205 the semi-ellipsoidal shell (an average length of  $a = 7.08 \pm 6.62$  mm, width of b = 5.25206  $\pm$  0.49 mm, and height of  $c = 3.14 \pm 0.54$  mm) formed by the two elytra into a projected 207 circle. We subdivided the projected circle into three ring-like regions, namely the inner, 208 central, and outer rings, to quantify the landing region (figure 3). The inner, central, and outer rings are in concordance with the landing angles of 0°-30°, 30°-60°, 60°-90°, 209 210 respectively.



**Figure 3.** Diagram of impact regions. The inner ring, central ring, and outer ring on the elytra shell are shown in different colors. The probabilities of the ladybird landing on inner ring, central ring, and outer ring are 20%, 20%, and 60%, respectively.

After dropping 20 insects for 60 times, the probability of collisions in different regions of the elytra was measured to be 20%, 20%, and 60% for the inner, central, and outer rings, respectively (Supplementary Video S1). Figures 4a-4c show three typical landing orientations with the landing angles of  $0^{\circ}$ ,  $30^{\circ}$ , and  $60^{\circ}$ . We also noticed that the elytra may open up and split on account of impact during landing on the inner ring (figure 4(d)). To quantify the opening of the elytra after a collision, we measured the splitting angle, which was defined as the angle between the elytra. Consequently, we categorized the fall into four cases: landing with closed elytra with collision in the inner ring (case 1), central ring (case 2), and outer ring (case 3), and landing with elytra splitting with collision in the inner ring (case 4).



**Figure 4.** Snapshots of the falling insects. (a)-(c) Snapshots showing the different landing angles. (d) Snapshots showing elytra splitting under impact.

## 221 3.2. Kinematics of falls

222 To quantify the kinematics of falls, we measured the height, translational velocities 223 and angular velocities of the ladybirds at 10 ms before and at 10 ms after collisions. For 224 the cases in which the elytra showed no splitting (case 1-3), the heights were 29.17  $\pm$ 225  $0.51 \text{ mm}, 28.65 \pm 0.72 \text{ mm}, \text{ and } 29.49 \pm 0.56 \text{ mm}, \text{ respectively (figure 5(a)), and the}$ 226 translational velocities were 2.79  $\pm$  0.09 m/s, 2.59  $\pm$  0.09 m/s, and 2.76  $\pm$  0.04 m/s 227 (figure 5(b)). Both heights and velocities vary only slightly, which indicates that 228 kinematics of falls are almost consistent between case 1-3, prior to collisions. In 229 contrast, we found that the heights and velocities in the rebound were noticeably 230 different. The rebound height was smaller in larger landing angles. Specifically, the 231 rebound height of the ladybirds when they landed on the inner ring (landing angle:  $0^{\circ}$ -232  $30^{\circ}$ ), was  $15.41 \pm 1.01$  mm, which is only 3.08% of the initial height. This is about 233 twice the rebound height of the ladybirds that landed on the outer ring (landing angle: 234  $60^{\circ}$ -90°). The translational velocity in the rebound also decreased from a peak value of 235  $1.06 \pm 0.08$  m/s when the ladybirds landed on the inner ring to  $0.31 \pm 0.06$  m/s when 236 they landed on the outer ring. After landing, the ladybird bodies rapidly turned at an 237 angular velocity of  $540.87 \pm 24.84$  rad/s when they landed on the outer ring, whereas 238 the angular velocities of the ladybirds landing on the central and inner rings were 239 smaller and equal to  $427.50 \pm 14.72$  rad/s, and  $305.25 \pm 29.02$  rad/s, respectively. The 240 observed rotations are likely due to the eccentric impact force generated at landing.

241 We found that the mechanical energy absorbed after landing were 68.46%, 57.46%, 242 and 59.15% for collisions on the inner, central, and the outer rings, respectively. 243 According to this, a ladybird landing on the inner ring is likely to absorb more impact 244 energy than others (figure 5(c)). Moreover, force plate recorded the peak forces as 37.09 245  $\pm$  1.51 mN, 30.23  $\pm$  0.37 mN, and 17.29  $\pm$  0.71 mN for impacts on the inner, central, 246 and outer rings, respectively (figure 5(d)). Compared to landing on the outer ring, the 247 peak impact force exerted on the inner ring was 214.52% larger. This inserts a force 248 that is 14 times the ladybird body weight, and is likely to be the cause for the elytra 249 splitting. However, interestingly, when splitting took place upon the collision, the 250 rebound height, translational velocity, and angular velocity decreased to  $9.09 \pm 1.22$ 251 mm,  $0.44 \pm 0.02$  m/s, and  $139.50 \pm 12.38$  rad/s, respectively, compared to the closed 252 elytra landing on the inner ring. Taking this into account, the energy absorption with 253 the elytra splitting reached 93.16% of the initial mechanical energy. This suggests that 254 the elytra splitting can dissipate the energy that could otherwise damage the ladybird 255 body. The elytra splitting also reduced the impact force to  $26.41 \pm 1.76$  mN, which is 256 71.21% of that generated in collisions when the elytra remain closed.



**Figure 5.** Kinematics of falls and impacts. (a) Height of the ladybirds from the ground. (b) Angular and translational velocities of the ladybirds before and after landing. (c) Energy absorption by the elytra during landing. (d) Peak impact force of the ladybird landing.

#### 257 3.3. The role of the morphology and structure on the energy absorption in falls

To understand how the elytra splitting can increase the energy absorption in collisions, we developed a detailed model of the ladybird body, as shown in figure 6(a). The dome-like elytra are multi-layers structures, with an average thickness of  $39.63 \pm$ 1.01 µm. The elytra are coupled to each other by the tenon ( $35.89 \pm 1.61$  µm) that inserts into the mortise ( $33.84 \pm 0.97$  µm) (figure 6(b)).



**Figure 6.** Morphology and microstructure of the ladybird elytra. (a) Micro-CT scans of a ladybird. (b) Scanning electron microscopy images of the elytra showing the structure of coupling.

263 Based on the geometry of the elytra, we proposed two comparative models to 264 investigate the role of coupling in energy absorption during collisions, as shown in 265 figure 7(a). A continuum model was presented as a quarter of a spherical shell with a 266 radius R of 1.5 mm and a thickness d of 40  $\mu$ m. In contrast to that, we also developed a 267 split model that was identical to the continuum model, except that it was subdivided 268 into two equally sized parts, which were connected using a coupling structure. The 269 coupling structure was a simplified version of the real coupling, and had a semicircle 270 cross section with a radius of 20  $\mu$ m (figure 7(a)).

271 Here, we compared the mechanical response of the models at the landing angle of 272  $0^{\circ}$  (figure 7(b), Supplementary Video S2). The numerical simulation of collision 273 includes three stages, as falling (1-2 in figure 7(b)), landing (2-3 in figure 7(b)), and 274 rebound (3-4 in figure 7(b)). For the continuum model, the changes of internal energy 275 can be subdivided into four phases  $(P_1-P_4 \text{ in figure 7(c)})$ . Before the model contacts the 276 ground, the internal energy of the model remains zero  $(P_1)$ . During landing, the model 277 deforms and the internal energy increases to 1.64  $\mu$ J (P<sub>2</sub>). With the recovery of the 278 elastic deformation, the model rebounds and its internal energy gradually declines to 279 0.07  $\mu$ J in P<sub>3</sub>. Considering that the influences of the air resistance and structural damping were not considered in our simulations, in  $P_3$ , the model vibrates with a 280 frequency of  $7.48 \times 10^5$  rad/s. 281

282 The changes of the internal energy of the split model can be subdivided into five phases  $(Q_1-Q_5 \text{ in figure 7(c)})$ . The internal energy of the split model in  $Q_1$  and  $Q_2$  are 283 284 the same as those of the continuum model in  $P_1$  and  $P_2$ , respectively, indicating that two 285 models have the same state in these two phases. However, the split model experienced 286 a larger deformation in phase  $Q_3$ , causing that the internal energy of the model to 287 increase to 1.99  $\mu$ J. Hence, we suggest that the split model has an increased capacity to 288 absorb energy, in comparison to the continuum model, because of splitting. 289 Subsequently, with the release of elastic energy, the model also rebounds off the ground in phase  $Q_4$ , and vibrates with a frequency of  $1.83 \times 10^5$  rad/s ( $Q_5$ ). According to the 290 291 free vibration frequency of the models after rebounding off the ground, we used the 292 equation  $K = \omega^2 m$  to measure the structural stiffness of the models. We discovered that the stiffness of the continuum model is about 16.75 times that of the split model, which can explain the larger deformation of split model upon collision. A larger deformation induced by structural specificity, especially satisfying a range of specific physiological demands, can be found in other organs of insects, such as dragonfly wings (Rajabi *et al.*, 2015), and mosquito antennas (Saltin *et al.*, 2019).

The peak impact force of the split model is 66.67% of that of the continuum model (figure 7(d)). This means that elytra splitting can reduce the risk of damage in collisions. The increase in absorbed energy in the split model also reduces the rebound velocity (figure 7(e)). The rebound velocity of the split model is only half of that of the continuum model (1.35 m/s vs. 2.62m/s). This shows that the elytra splitting also reduces the kinetic energy of a falling ladybird during rebounding.

Besides its positive impacts, the elytra splitting may also have negative effects. The increased internal energy of the split model is accompanied by an increased level of stress within the model. As seen in figure 7(f), the peak average stress of the split model is 2.18 times that of the continuum model, which is caused by the comparatively larger deformation of the model. However, it seems that the stress developed within the split model is still too small to result in the failure of the material forming the elytra, which has a strength of about 72 MPa (Rajabi *et al.*, 2017).



**Figure 7.** Numerical simulation of the mechanical behavior of the elytra in collisions. (a) The continuum and split models. *R*: radius of the spherical shell; *d*: the thickness of the shell; *r*: radius of the coupling. (b) Models reach the ground with a landing angle of 0°. Simulation of falls include three stages, as falling (1-2), landing (2-3), and rebound (3-4). (c) Comparison of the internal energy of the models. Falling of the continuum includes four phases. *P*<sub>1</sub>: falling; *P*<sub>2</sub>: elastic deformation during landing; *P*<sub>3</sub>: shape recovery during rebound; *P*<sub>4</sub>: free vibration. Falling of the split model includes five phases: *Q*<sub>1</sub>: falling; *Q*<sub>2</sub>: first elastic deformation during landing; *Q*<sub>3</sub>: second elastic deformation during landing *Q*<sub>4</sub>: shape recovery during rebound; *Q*<sub>5</sub>: free vibration. (d)-(f) Comparison of the impact force, translational velocity and average stress of the models.

We further simulated the falls of the models with other landing angles  $(30^{\circ} \text{ and} 60^{\circ})$  (figures 8(a)-(b)). The split model always exhibited a higher energy absorption capacity, which reduced both the peak impact force (at least by 25.36%) and rebound velocity (at least by 3.40%), as shown in Table 2. The peak average stress of the split 315 model was at most 1.2 times larger than the other model. These results are consistent 316 with those obtained from the simulation of falls at  $0^{\circ}$  landing angle (figures 7(c)-(f)). 317 Additionally, we also found that with the increase of the landing angle, both the peak impact force and rebound velocity of the model reduced without absorbing more energy. 318 319 This means that the risk damage can be reduced by simply adjusting the landing angles without absorbing more impact energy. This may explain why the ladybirds often 320 321 landed on the outer ring, although it remains unclear whether this adjustment is done 322 by means of active or passive mechanisms.



Figure 8. Numerical simulation of the falls in the landing angles of 30° and 60°.

<sup>323</sup> 324

Table 2. Numerical simulation results of the models					
Model type	Landing angle	Internal	Peak impact	Rebound	Peak average
	(deg)	energy (µJ)	force (mN)	velocity (m/s)	stress (MPa)
Continuum	30	1.13	73.71	2.47	0.69
model	60	0.97	47.48	2.06	0.73
California dal	30	1.70	55.02	2.24	0.83
Spiit model	60	1.63	41.09	1.99	0.77

#### 325 **3.4.** Application

To demonstrate the application of the energy absorption mechanism discovered in this study, we fabricated the comparative models by 3D printing. We dropped the models from a height of 2 cm without initial velocity, and measured the impact forces

329	generated in collisions by a force plate (figure $9(a)$ ). The short dropping distance of 2
330	cm enabled us to always have the landing angle of $0^\circ$ in the experiments. The impact
331	force of the continuum model is $1.72\pm0.15$ N, whereas that of the split model is $72.67\%$
332	of the former (1.25 $\pm$ 0.10 N) (figure 9(b), Supplementary Video S3). To test the
333	performance of the insect-inspired 3D printed models in real-world applications, we
334	attached them as landers to the arms of a quadcopter (JY03, BBS, China) with a weight
335	of ~0.2 kg (figure 9(c)). We let the quadcopter fly and then turned it off in a height of
336	26 cm above the ground to simulate a free-fall scenario (figures 9(d)-(e),
337	Supplementary Video S4). The selected height was selected as it enabled us to capture
338	the whole falling event within the frame of the camera. Similar to the results obtained
339	from testing of isolated models, the peak impact force of the quadcopter equipped with
340	3D printed split models was smaller than, and only 74.85% of, that of the quadcopter
341	equipped with 3D printed continuum models (figure 9(f)). The quadcopter equipped
342	with 3D printed split models also experienced a smaller rebound height, in comparison
343	to the quadcopter equipped with 3D printed continuum models (6.0 cm vs. 6.4 cm).
344	As a result, it also became stable in a shorter time than the quadcopter equipped with
345	3D printed continuum models (1.14 s vs. 1.67 s).



**Figure 9.** Insect-inspired 3D printed models in application. (a) Experimental setup. (b) The impact force of the split model (with coupling) and continuum model (without coupling) in falls. (c) A quadcopter equipped with four insect-inspired 3D printed models. (d)-(e) Comparison of the mechanical performance of the 3D printed split models and the 3D continuum models attached to the arms of a quadcopter. (f) Comparison of peak impact forces.

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## 347 **4. Conclusion**

348 The results of our study suggest that elytra coupling may act as a mechanism to 349 increase the energy absorption in intentional falls after encountering predators. This mechanism is achieved by the failure of the coupling upon excessive loads and may 350 351 also exist in other beetle species. Three directions for future research seem particularly 352 worth following. First, the intentional falls on the elytra may be achieved by the combined attitude regulation and in-flight maneuver (Zheng et al., 2017), but there 353 354 lacks extensive investigations on this problem. Second, comparative study of the 355 mechanical behavior of the coupling mechanisms in insects that show a similar behavior can shed light on the diversity of this strategy. Third, detailed investigation of the 356 357 specialized morphology and material properties of the coupling mechanism may inspire

358	the biomimetic design of insect-inspired energy absorption devices.
359	
360	Competing interests
361	We declare we have no competing interests.
362	
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371	
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